

90(2) : Complete Solution for Precessing Orbits

In case it is assumed that the scalar potential is, or an excellent approximation:

$$\Phi = -\frac{MG}{(x^2 + y^2)^{1/2}}, \quad (1)$$

because precessional effects are small.

It follows as in Note 389(4) that for forward precession:

$$\omega_x = \frac{x}{x^2 + y^2} \left(\frac{1}{y} - 1 \right) - \frac{\dot{x}\dot{y}y + x\dot{x}^2}{\gamma c^2 (x^2 + y^2)} \quad (2)$$

$$\omega_y = \frac{y}{x^2 + y^2} \left(\frac{1}{x} - 1 \right) - \frac{\dot{y}\dot{x}x + y\dot{y}^2}{\gamma c^2 (x^2 + y^2)} \quad (3)$$

and for retrograde precession:

$$\omega_x = \left(\frac{x}{x^2 + y^2} \right) \left(\left(1 - \frac{\dot{x}^2 + \dot{y}^2}{c^2} \right)^{3/2} - 1 \right) \quad (4)$$

$$\omega_y = \left(\frac{y}{x^2 + y^2} \right) \left(\left(1 - \frac{\dot{x}^2 + \dot{y}^2}{c^2} \right)^{3/2} - 1 \right) \quad (5)$$

$$\text{where } \gamma = \left(1 - \frac{\dot{x}^2 + \dot{y}^2}{c^2} \right)^{-1/2} \quad (6)$$

The spin corrections have been found from conservation of scalar antisymmetry.

From conservation of vector antisymmetry:

$$\frac{\partial Q_z}{\partial y} + \frac{\partial Q_y}{\partial z} = \omega_y Q_z + \omega_z Q_y \quad (7)$$

$$\frac{\partial Q_x}{\partial z} + \frac{\partial Q_z}{\partial x} = \omega_z Q_x + \omega_x Q_z \quad (8)$$

$$\frac{\partial Q_y}{\partial x} + \frac{\partial Q_x}{\partial y} = \omega_x Q_y + \omega_y Q_x \quad (9)$$

These have particular solutions:

$$\left. \begin{aligned} \frac{\partial Q_z}{\partial y} &= \omega_y Q_z, \quad \frac{\partial Q_y}{\partial z} = \omega_z Q_y, \\ \frac{\partial Q_x}{\partial z} &= \omega_z Q_x, \quad \frac{\partial Q_z}{\partial x} = \omega_x Q_z, \\ \frac{\partial Q_y}{\partial x} &= \omega_x Q_y, \quad \frac{\partial Q_x}{\partial y} = \omega_y Q_x \end{aligned} \right\} - (10)$$

$$\omega_z = 0. - (11)$$

in which:

Therefore:

$$\left. \begin{aligned} \frac{\partial Q_z}{\partial y} &= \omega_y Q_z, \quad \frac{\partial Q_y}{\partial z} = 0, \\ \frac{\partial Q_x}{\partial z} &= 0, \quad \frac{\partial Q_z}{\partial x} = \omega_x Q_z, \\ \frac{\partial Q_y}{\partial x} &= \omega_x Q_y, \quad \frac{\partial Q_x}{\partial y} = \omega_y Q_x \end{aligned} \right\} - (12)$$

The general solution of the first order differential eqn.:

$$\frac{dy}{dx} + y p(x) = q(x) - (13)$$

$$y = \frac{1}{u(x)} \left(\int u(x) q(x) dx + C \right) - (14)$$

is

$$u(x) = \exp \left(\int p(x) dx \right) - (15)$$

where

$$q(x) = 0 - (16)$$

If:

$$\frac{dy}{dx} = -y p(x) - (17)$$

then

whose solution is:

$$y = C \exp \left(- \int p(x) dx \right) - (18)$$

$$Q_x = Q_x(0) \exp \left(\int \omega_y dy \right) - (19)$$

$$Q_y = Q_y(0) \exp \left(\int \omega_x dx \right) - (20)$$

$$Q_z = Q_z(0) \exp \left(\int \omega_y dy \right) - (21)$$

$$= Q_z(0) \exp \left(\int \omega_x dx \right)$$

which:

$$\frac{\partial Q_x}{\partial z} = \frac{\partial Q_y}{\partial z} = 0 - (22)$$

The vector potential is:

$$\underline{Q} = Q_x \underline{i} + Q_y \underline{j} + Q_z \underline{k} - (23)$$

general law:

$$\int \omega_y dy \neq \int \omega_x dx - (24)$$

from eq. (21):

$$Q_z(0) = 0 - (25)$$

only self consistent solution. So:

$$Q_z = 0 - (26)$$

$$\underline{Q} = Q_x \underline{i} + Q_y \underline{j} - (27)$$

The gauge magnetic field is found from \underline{Q} and

$$\underline{\underline{Q}} = \underline{\nabla} \times \underline{\underline{Q}} - \underline{\omega} \times \underline{\underline{Q}} - (28)$$

The scalar component of the spin connection is found from conservation of trace antisymmetry (the Lichnerowicz constraint):

$$\frac{1}{c^2} \left(\frac{\partial}{\partial t} + \omega_0 \right) \underline{\Phi} = (\underline{\nabla} - \underline{\omega}) \cdot \underline{\underline{Q}} - (29)$$

From eq. (1): $\frac{\partial \underline{\Phi}}{\partial t} = 0 - (30)$

so $\omega_0 = \frac{c^2}{\underline{\Phi}} (\underline{\nabla} - \underline{\omega}) \cdot \underline{\underline{Q}} - (31)$

The scalar spin connection must also satisfy the scalar law of conservation of antisymmetry:

$$\underline{g} = -\underline{\nabla} \underline{\Phi} + \underline{\omega} \underline{\Phi} = -\frac{\partial \underline{\underline{Q}}}{\partial t} - \omega_0 \underline{\underline{Q}} - (32)$$

If it is assumed that:

$$\frac{\partial \underline{\underline{Q}}}{\partial t} = \underline{0} - (33)$$

then $\underline{g} = -\omega_0 \underline{\underline{Q}} - (34)$

i.e. $(g_x^2 + g_y^2)^{1/2} = -\omega_0 (Q_x^2 + Q_y^2)^{1/2} - (35)$

If ω_0 from eq. (35) is not the same as ω_0

from eq (31), the regauge $\underline{\Phi}$ and \underline{Q} :

$$\underline{\Phi} \rightarrow \underline{\Phi} + \frac{\partial \phi}{\partial t} \quad - (36)$$

where ϕ is a gauge function.

The method of regauging \underline{Q} is to define:

$$\underline{g} := -\underline{\nabla} \underline{\Phi} - \frac{\partial \underline{Q}}{\partial t} \text{ (total)} \quad - (37)$$

$$\underline{\Omega} = \underline{\nabla} \times \underline{Q} \text{ (total)} \quad - (38)$$

where

$$\underline{Q} \text{ (total)} = \underline{Q} + \underline{Q}_1 \quad - (39)$$

and

$$\underline{\nabla} \times \underline{Q}_1 = -\underline{\omega} \times \underline{Q} \quad - (40)$$

The four potential to be regauged is defined as:

$$Q^4 \text{ (total)} := \left(\frac{\underline{\Phi}}{c}, \underline{Q} \text{ total} \right) \quad - (41)$$

and the regauging is:

$$\underline{\Phi} \rightarrow \underline{\Phi} + \frac{\partial \phi}{\partial t} \quad - (42)$$

$$\underline{Q} \text{ total} \rightarrow \underline{Q} \text{ total} - \underline{\nabla} \phi \quad - (43)$$

under which the observable \underline{g} and $\underline{\Omega}$ are unchanged.

Eqs. (37) and (38) satisfy the ECE2 homogeneous field equations of gravitation:

$$\underline{\nabla} \cdot \underline{\Omega} = 0 \quad - (44)$$

$$\underline{\nabla} \times \underline{g} + \frac{\partial \underline{\Omega}}{\partial t} = \underline{0}$$

By definition:

$$-\frac{dQ(\text{total})}{dt} = \underline{\omega} \underline{\Phi} \quad - (45)$$

Regarding this equation means that: - (46)

$$-\frac{d}{dt} (Q(\text{total}) - \nabla \phi) = \underline{\omega} \left(\underline{\Phi} + \frac{d\phi}{dt} \right)$$

in eqs. (45) and (46):

$$\frac{d}{dt} \nabla \phi = \underline{\omega} \frac{d\phi}{dt} \quad - (47)$$

i.e.

$$\nabla \frac{d\phi}{dt} = \underline{\omega} \frac{d\phi}{dt} \quad - (48)$$

Also, note that:

$$Q(\text{total}) = - \int \underline{\omega} \underline{\Phi} dt = - \underline{\omega} \int \underline{\Phi} dt \quad - (49)$$

If

$$\underline{\Phi} \rightarrow \underline{\Phi} + \frac{d\phi}{dt} \quad - (50)$$

then

$$Q(\text{total}) \rightarrow - \int \underline{\omega} \underline{\Phi} dt - \phi \underline{\omega} \quad - (51)$$

$$= Q(\text{total}) - \nabla \phi$$

so

$$\boxed{\nabla \phi = \underline{\omega} \phi} \quad - (52)$$

Regarding means that:

$$Q(\text{total}) \rightarrow Q(\text{total}) - \underline{\omega} \phi \quad - (53)$$

$$\underline{\Phi} \rightarrow \underline{\Phi} + \frac{d\phi}{dt} \quad - (54)$$

By definition:

$$\underline{Q}(\text{total}) = \underline{Q} + \underline{Q}_1 - (55)$$

2nd

$$\underline{\nabla} \times \underline{Q}_1 = -\underline{\omega} \times \underline{Q} - (56)$$

3rd

$$\underline{Q} + \underline{Q}_1 \rightarrow \underline{Q} + \underline{Q}_1 - \underline{\omega} \phi - (57)$$

Since ϕ is arbitrary, it can always be assumed that:

$$\underline{Q} \rightarrow \underline{Q} - \underline{\omega} \phi - (58)$$

This means that:

$$\begin{aligned} \underline{\nabla} \times \underline{Q}_1 &\rightarrow \underline{\nabla} \times \underline{Q}_1 - \underline{\omega} \times \underline{\omega} \phi - (59) \\ &= \underline{\nabla} \times \underline{Q}_1 \end{aligned}$$

Also:

$$\underline{Q}_1 \rightarrow \underline{Q}_1 - \underline{\omega} \phi - (60)$$

Therefore the gauge transformation is:

$\underline{\Phi} \rightarrow \underline{\Phi} + \frac{d\phi}{dt} - (61)$
$\underline{Q} \rightarrow \underline{Q} - \underline{\omega} \phi - (62)$
$\underline{\nabla} \phi = \underline{\omega} \phi - (63)$

From eq. (63):

$$\frac{d\phi}{dx} = \omega_x \phi - (64)$$

$$\frac{d\phi}{dy} = \omega_y \phi - (65)$$

so

$$\psi = \psi_1(0) \exp \left(\int \omega_x dx \right) = \psi_2(0) \exp \left(\int \omega_y dy \right) \quad - (66)$$

Here $\psi_1(0)$ and $\psi_2(0)$ are constants of integration.

From eq. (62):

$$Q_x \rightarrow Q_x - \omega_x \psi \quad - (67)$$

$$Q_y \rightarrow Q_y - \omega_y \psi \quad - (68)$$

so eq. (35) is reorganized to:

$$\begin{aligned} (g_x^2 + g_y^2)^{1/2} &= -\omega_0 (Q_x^2 + Q_y^2)^{1/2} \\ &\rightarrow -\omega_0 \left((Q_x - \omega_x \psi)^2 + (Q_y - \omega_y \psi)^2 \right)^{1/2} \end{aligned} \quad - (69)$$

and ψ is adjusted so that ω_0 from eq. (69) is the same as ω_0 from eq. (31).

Finally, for eq. (2) of Note 390(1), the following equation must be solved:

$$\nabla \cdot (\underline{0} \quad \underline{\Phi}) = 0 \quad - (70)$$

This is solved in the Newtonian limit, where:

$$\omega_x = \omega_y \rightarrow 0 \quad - (71)$$

in eqs. (2) and (3). More generally, eqs. (1) to (3) and (70) must be solved numerically.