

### 352(3): Time Dependent Equations

These are:  $\frac{\partial \underline{v}}{\partial t} + \underline{\nabla} h = -(\underline{v} \cdot \underline{\nabla}) \underline{v} \quad - (1)$

$$\frac{\partial h}{\partial t} + a_0^2 \underline{\nabla} \cdot \underline{v} = -\underline{v} \cdot \underline{\nabla} h \quad - (2)$$

$$\frac{\partial \underline{w}}{\partial t} + \underline{\nabla} \times (\underline{w} \times \underline{v}) = \frac{1}{R} \nabla^2 \underline{w} \quad - (3)$$

$$\frac{\partial \underline{w}}{\partial t} + \underline{\nabla} \times \underline{E} = \underline{0} \quad - (4)$$

$$\underline{E} = -\frac{\partial \underline{v}}{\partial t} - \underline{\nabla} h \quad - (5)$$

$$a_0^2 \underline{\nabla} \times \underline{H} - \frac{\partial \underline{E}}{\partial t} = \underline{J} \quad - (6)$$

$$\underline{H} = \underline{w} \quad - (7)$$

where

$$\underline{J} = \frac{\partial^2 \underline{v}}{\partial t^2} + \underline{\nabla} \frac{\partial h}{\partial t} + a_0^2 \underline{\nabla} \times (\underline{\nabla} \times \underline{v}) \quad - (8)$$

and

$$\underline{\nabla} h = \frac{1}{\rho} \Delta \underline{p} \quad - (9)$$

From eqs. (4) and (5):

$$\frac{\partial \underline{w}}{\partial t} - \frac{\partial \underline{\nabla} \times \underline{v}}{\partial t} - \underline{\nabla} \times \underline{\nabla} h = \underline{0} \quad - (10)$$

i.e.  $\frac{\partial \underline{w}}{\partial t} = \underline{\nabla} \times \frac{\partial \underline{v}}{\partial t} \quad - (11)$

which is consistent with:

$$\underline{w} = \underline{\nabla} \times \underline{v} \quad - (12)$$

Also from eqs. (4) and (5):

$$2) \quad \frac{\partial \underline{w}}{\partial t} + \underline{\nabla} \times (\underline{v} \cdot \underline{\nabla}) \underline{v} = \underline{0} \quad - (13)$$

From eqs. (12) and (13):

$$\boxed{\frac{D \underline{v}}{Dt} = \frac{\partial \underline{v}}{\partial t} + (\underline{v} \cdot \underline{\nabla}) \underline{v} = \underline{0}} \quad - (14)$$

is a particular solution. Eq. (14) means that:

$$\underline{\nabla} h = \frac{1}{\rho} \Delta p = \underline{0} \quad - (15)$$

and there is no change of pressure.

Also from eq. (4):

$$\underline{\nabla} \times \frac{\partial \underline{v}}{\partial t} + \underline{\nabla} \times \underline{E} = \underline{0}, \quad - (16)$$

a particular solution of  $\frac{\partial \underline{v}}{\partial t}$  which is:

$$\frac{\partial \underline{v}}{\partial t} + \underline{E} = \underline{0}. \quad - (17)$$

$$\text{If} \quad \underline{E} = - \frac{\partial \underline{v}}{\partial t} - \underline{\nabla} h, \quad - (18)$$

eqs. (17) and (18) imply:

$$\underline{\nabla} h = \underline{0} \quad - (19)$$

again.

The simple flow of Eq. (14) can be graphed and animated as a preliminary exercise.

The Beltrami solution of Eq. (3) is

$$\boxed{\frac{\partial \underline{w}}{\partial t} = \frac{1}{R} \nabla^2 \underline{w}} \quad - (20)$$

because for Beltrami flow:

$$\underline{w} \times \underline{v} = \underline{0} \quad - (21)$$

and

$$\underline{\nabla} \times \underline{v} = R \underline{v} \quad - (22)$$

The simple equation (20) can also be graphed and animated. At a certain Reynolds number the vorticity of Beltrami flow develops turbulence.

In general, eq. (3) is:

$$\begin{aligned} \frac{d}{dt} (\underline{\nabla} \times \underline{v}) + \underline{\nabla} \times (\underline{w} \times \underline{v}) &= \frac{1}{R} \underline{\nabla}^2 \underline{w} \\ &= \frac{1}{R} (\underline{\nabla} (\underline{\nabla} \cdot \underline{w}) - \underline{\nabla} \times (\underline{\nabla} \times \underline{w})) \end{aligned} \quad - (23)$$

in which  $\underline{\nabla} \cdot \underline{w} = \underline{\nabla} \cdot \underline{\nabla} \times \underline{v} = 0 \quad - (24)$

Therefore a particular solution of eq. (23) is:

$$\boxed{\frac{d\underline{v}}{dt} + \underline{w} \times \underline{v} = -\frac{1}{R} \underline{\nabla} \times \underline{w}} \quad - (25)$$

so  $\underline{\nabla} \times \underline{w} = -R \left( \frac{d\underline{v}}{dt} + \underline{w} \times \underline{v} \right) \quad - (26)$

For Beltrami flow:

$$\underline{\nabla} \times \underline{w} = -R \frac{d\underline{v}}{dt} \quad - (27)$$

4) Therefore Beltrami flow is described by eqs (20) and (27):

$$\nabla^2 \underline{w} = R \frac{\partial \underline{w}}{\partial t} \quad - (28)$$

and  $\nabla \times \underline{w} = -R \frac{\partial \underline{v}}{\partial t} \quad - (29)$

with  $\underline{w} = \nabla \times \underline{v} \quad - (30)$

Therefore  $\nabla^2 \underline{v} = R \frac{\partial \underline{v}}{\partial t} \quad - (31)$

and  $\nabla \times \underline{w} = -R \frac{\partial \underline{v}}{\partial t} \quad - (32)$

$$\begin{aligned} &= \nabla \times (\nabla \times \underline{v}) \\ &= \nabla (\nabla \cdot \underline{v}) - \nabla^2 \underline{v} \quad - (33) \end{aligned}$$

These equations imply that Beltrami flow is incompressible and a divid because:

$$\nabla \cdot \underline{v} = 0 \quad - (34)$$

The interesting features are obtained therefore when:

$$\underline{w} \times \underline{v} \neq 0 \quad - (35)$$

and  $\frac{\partial \underline{v}}{\partial t} + (\nabla \times \underline{v}) \times \underline{v} = -\frac{1}{R} \nabla \times (\nabla \times \underline{v}) \quad - (36)$

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