

### 51(5): Simplified Derivation of Relativistic Zeeman Splitting

Start with the relativistic classical Hamiltonian:

$$H_1 = -\frac{e}{2m} \left(1 - \frac{p_0^2}{m^2 c^2}\right)^{-1/2} \underline{L}_0 \cdot \underline{B} \quad - (1)$$

$$\sim -\frac{e}{2m} \left(1 + \frac{1}{2} \frac{p_0^2}{m^2 c^2}\right) \underline{L}_0 \cdot \underline{B}$$

Quantize by assuming that:

$$\hat{H}_1 \psi = -\frac{e}{2m} \left(1 + \frac{1}{2} \frac{p_0^2}{m^2 c^2}\right) \underline{B} \cdot \hat{\underline{L}}_0 \psi \quad - (2)$$

in which  $\hat{\underline{L}}_0$  is an operator and  $p_0^2$  a function. If  $\underline{B}$  is aligned in the  $z$  axis:

$$\hat{H}_1 \psi = -\frac{e}{2m} \left(1 + \frac{1}{2} \frac{p_0^2}{m^2 c^2}\right) B_z \hat{L}_{0z} \psi \quad - (3)$$

where

$$\hat{L}_{0z} \psi = m_L \hbar \psi \quad - (4)$$

with

$$m_L = -L, \dots, L \quad - (5)$$

So the energy levels are:

$$H_1 = \langle \hat{H}_1 \rangle = -\frac{e\hbar}{2m} \left(1 + \frac{1}{2} \frac{p_0^2}{m^2 c^2}\right) m_L \quad - (6)$$

Finally use:

$$\frac{p_0^2}{2m} = \left\langle \frac{\hat{p}_0^2}{2m} \right\rangle - (7)$$

So

$$H_1 = -\frac{e\hbar}{2m} m_L \left( 1 + \frac{1}{mc^2} \left\langle \frac{\hat{p}_0^2}{2m} \right\rangle \right) - (8)$$

$$= -\frac{e\hbar}{2m} m_L \left( 1 - \frac{\hbar^2}{4m^2 c^2} \int \psi^* \nabla^2 \psi d\tau \right)$$

In a rigorous development,  $\psi$  is the relativistic wave function of an atom or molecule.

In a rough first approximation, the non-relativistic wave functions of atomic H are used:

$$\left\langle \frac{\hat{p}_0^2}{2m} \right\rangle = \frac{m e^4}{32\pi^2 \epsilon_0^2 \hbar^2 n^3} - (9)$$

So:

$$H_1 = -\frac{e\hbar}{2m} m_L \left( 1 + \left( \frac{e^4}{32\pi^2 \epsilon_0^2 \hbar^2 c^2} \right) \frac{1}{n^3} \right) - (10)$$

This result can be expressed as:

$$H_1 = -\frac{e\hbar}{2m} m_L \left( 1 + \frac{1}{2} \left( \frac{\lambda_c}{a_0} \right) \frac{1}{n^3} \right) - (11)$$

where:

$$3) \lambda_c = \frac{h}{mc} = 3.861591 \times 10^{-13} \text{ m} - (12)$$

is the Compton wavelength of the electron,

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{me^2} = 5.29177 \times 10^{-11} \text{ m} - (13)$$

is the Bohr radius, and

$$\alpha = \frac{e^2}{4\pi\hbar c \epsilon_0} = 0.007297351 - (14)$$

is the fine structure constant.

So:

$$H_1 = \langle H_1 \rangle = -\frac{e\hbar}{2m} m_L B_z \left( 1 + \frac{2.662567 \times 10^{-5}}{n^2} \right) - (15)$$

Except for a sign change, this is the result used in Note 331(3) to demonstrate relativistic Zeeman splitting when the Dirac approximation is no longer used.

---