

# 18(5): Spa Orbit Interaction for Fermi Equation

We refer to:

P. Novak, "Calculation of Spn-Orbit Coupling"  
Institute of Physics, Czech Republic.

The fermi equation is:

$$(\hat{E} - c \underline{\sigma} \cdot \hat{\underline{p}}) \phi^R = mc^2 \phi^L \quad (1)$$

$$(\hat{E} + c \underline{\sigma} \cdot \hat{\underline{p}}) \phi^L = mc^2 \phi^R \quad (2)$$

In the presence of a potential  $V$ :

$$V = -e\phi \quad (3)$$

Then:

$$(\hat{E} + V - c \underline{\sigma} \cdot \hat{\underline{p}}) \phi^R = mc^2 \phi^L \quad (4)$$

$$(\hat{E} + V + c \underline{\sigma} \cdot \hat{\underline{p}}) \phi^L = mc^2 \phi^R \quad (5)$$

As in UFT 173 make the transformations:

$$\phi^R = \frac{1}{\sqrt{2}} (\phi_S^R + \phi_S^L) \quad (6)$$

$$\phi^L = \frac{1}{\sqrt{2}} (\phi_S^R - \phi_S^L) \quad (7)$$

to obtain:

$$(\hat{E} - V - mc^2) \phi_S^R = c \underline{\sigma} \cdot \hat{\underline{p}} \phi_S^L \quad (8)$$

$$(\hat{E} - V + mc^2) \phi_S^L = c \underline{\sigma} \cdot \hat{\underline{p}} \phi_S^R \quad (9)$$

Now write

$$\epsilon = \hat{E} - mc^2 \quad (10)$$

$$\phi_S^R = \underline{\Phi}, \quad \phi_S^L = \chi \quad (11)$$

to obtain:

$$c \underline{\sigma} \cdot \hat{\underline{p}} \chi = (\epsilon - V) \underline{\Phi} \quad (12)$$

$$c \underline{\sigma} \cdot \hat{\underline{p}} \underline{\Phi} = (\epsilon - V + 2mc^2) \chi \quad (13)$$

2) There are Novak's eqns. (4) and (5), which have already been derived from ECE theory

From eq. (13):

$$\chi = \frac{c \underline{\sigma} \cdot \underline{\hat{p}} \underline{\Phi}}{E + V + 2mc^2} \quad (14)$$

So in eq. (12):

$$E \underline{\Phi} = -V \underline{\Phi} + c^2 \underline{\sigma} \cdot \underline{\hat{p}} \left( \frac{\underline{\sigma} \cdot \underline{\hat{p}} \underline{\Phi}}{E + V + 2mc^2} \right) \quad (15)$$

$$= -V \underline{\Phi} + \frac{1}{2m} \underline{\sigma} \cdot \underline{\hat{p}} \left( \frac{\underline{\sigma} \cdot \underline{\hat{p}} \underline{\Phi}}{1 + \frac{E + V}{2mc^2}} \right) \quad (16)$$

This is Novak's eq. (6).

Now we:

$$\left( 1 + \frac{E + V}{2mc^2} \right)^{-1} \sim 1 - \left( \frac{E + V}{2mc^2} \right) \quad (17)$$

$$\text{if } E \neq V \ll 2mc^2 \quad (18)$$

Therefore:

$$E \underline{\Phi} = -V \underline{\Phi} + \frac{1}{2m} \underline{\sigma} \cdot \underline{\hat{p}} \left( \left( 1 - \frac{E + V}{2mc^2} \right) \underline{\sigma} \cdot \underline{\hat{p}} \underline{\Phi} \right) \quad (19)$$

where

$$\underline{\hat{p}} = -i \underline{\nabla} \quad (20)$$

3) Eq. (19) is:

$$\hat{E}\Phi = -\nabla^2\Phi + \frac{1}{2m}\underline{\sigma}\cdot\underline{\hat{p}}\underline{\sigma}\cdot\underline{\hat{p}}\Phi - \frac{1}{4m^2c^2}\underline{\sigma}\cdot\underline{\hat{p}}\left(\hat{E}\underline{\sigma}\cdot\underline{\hat{p}}\Phi\right) + \frac{1}{4m^2c^2}\underline{\sigma}\cdot\underline{\hat{p}}\left(\nabla\underline{\sigma}\cdot\underline{\hat{p}}\Phi\right) \quad (21)$$

Norvik makes the non-relativistic approximation:

$$T = E = E - mc^2 = (\gamma - 1)mc^2 \quad (22)$$

$$\rightarrow \frac{p^2}{2m}$$

$$v \ll c. \quad (23)$$

for

So:

$$\hat{E}\Phi = -\nabla^2\Phi + \frac{p^2}{2m}\Phi - \frac{p^4}{8m^3c^2}\Phi + \frac{1}{4m^2c^2}\underline{\sigma}\cdot\underline{\hat{p}}\left(\nabla\underline{\sigma}\cdot\underline{\hat{p}}\Phi\right) \quad (24)$$

where we have used:

$$\underline{\sigma}\cdot\underline{\hat{p}}\underline{\sigma}\cdot\underline{\hat{p}} = p^2 \quad (25)$$

The Schrodinger equation is the approximation:

$$\hat{E}\Phi = \left(\nabla^2 + \frac{p^2}{2m}\right)\Phi = \hat{H}\Phi \quad (26)$$

The third term on the RHS of eq. (24) is known as the mass term.

4) The first term on the RHS of eq. (24) gives the spin-orbit interaction term and the Darwin term.

This term is:

$$\frac{\underline{\sigma} \cdot \underline{\hat{p}}}{4m^2 c^2} (\underline{\nabla} \underline{\sigma} \cdot \underline{\hat{p}} \Phi) = -\frac{i\hbar}{4m^2 c^2} \underline{\sigma} \cdot \underline{\nabla} (\underline{\nabla} \underline{\sigma} \cdot \underline{\hat{p}} \Phi) \quad - (27)$$

The RHS is expanded using the Leibniz Theorem, so

$$\text{RHS} = -\frac{i\hbar}{4m^2 c^2} \left( \underline{\sigma} \cdot \underline{\nabla} \underline{\nabla} \underline{\sigma} \cdot \underline{\hat{p}} \Phi + \underline{\nabla} (\underline{\sigma} \cdot \underline{\nabla}) (\underline{\sigma} \cdot \underline{\hat{p}} \Phi) \right) \quad - (28)$$

The electric field is defined as:

$$\underline{V} = e\phi; \quad e\underline{E} = -\underline{\nabla} \underline{V}, \quad - (29)$$

so:

$$\text{RHS} = -\frac{i\hbar e}{4m^2 c^2} \left( \underline{\sigma} \cdot \underline{E} \underline{\sigma} \cdot \underline{\hat{p}} \Phi + \underline{\nabla} (\underline{\sigma} \cdot \underline{\nabla}) (\underline{\sigma} \cdot \underline{\hat{p}} \Phi) \right) \quad - (30)$$

By Pauli algebra:

$$\underline{\sigma} \cdot \underline{E} \underline{\sigma} \cdot \underline{\hat{p}} = \underline{E} \cdot \underline{\hat{p}} + i \underline{\sigma} \cdot \underline{E} \times \underline{\hat{p}} \quad - (31)$$

Therefore:

$$\begin{aligned} \text{RHS} = & -\frac{e\hbar}{4m^2 c^2} \underline{\sigma} \cdot \underline{E} \times \underline{\hat{p}} - \frac{\hbar^2}{4m^2 c^2} \underline{\nabla} \underline{\nabla} \underline{\nabla} \Phi \\ & - \frac{\hbar^2}{4m^2 c^2} \underline{\nabla} \underline{\nabla}^2 \Phi \quad - (32) \end{aligned}$$

5) Eq. (21) is therefore:

$$\hat{H} \Phi = E \Phi \quad - (33)$$

where:

$$\hat{H} = e\phi + \frac{\hat{p}^2}{2m} - \frac{\hat{p}^4}{8m^3c^2} - \frac{e\hbar}{4m^2c^2} \underline{\sigma} \cdot \underline{E} \times \underline{\hat{p}} - \frac{\hbar^2}{4m^2c^2} \underline{\nabla} \underline{\nabla} \underline{\nabla} - \frac{\hbar^2 \underline{\nabla} \underline{\nabla}^2}{4m^2c^2} \quad - (34)$$

The first term is the spin-orbit coupling term, the fifth term is the Darwin term, and the sixth term is a relativistic correction of  $\hat{p}^2/(2m)$ .

These terms all emerge from the fermion equation.

In a Coulomb potential:

$$\phi = -\frac{e}{4\pi\epsilon_0 r} \quad - (35)$$

$$\underline{E} = -\frac{e^2}{4\pi\epsilon_0} \frac{\underline{r}}{r^3} \quad - (36)$$

so the spin orbit term is:

$$\hat{H}_{so} = -\frac{e^2 \hbar}{16\pi\epsilon_0 r^3 m^2 c^2} \underline{\sigma} \cdot \underline{r} \times \underline{\hat{p}} \quad - (37)$$

$$= -\frac{e^2 \hbar}{16\pi\epsilon_0 r^3 m^2 c^2} \underline{\sigma} \cdot \underline{L} \quad - (38)$$

6) where  $\hat{\underline{L}}$  is the orbital angular momentum operator.  
The spin angular momentum is defined as an operator

$$\hat{\underline{S}} = \frac{1}{2} \hbar \underline{\underline{\sigma}} \quad - (39)$$

So:

$$\hat{H}_{so} = + \frac{e^2}{8\pi\epsilon_0 r^3 m^2 c^2} \hat{\underline{S}} \cdot \hat{\underline{L}} \quad - (40)$$

Therefore the relativistic Hamiltonian (34) is:

$$\hat{H} = \nabla^2 - \frac{\hbar^2}{2m} \left( 1 + \frac{\nabla^2}{2m^2 c^2} \right) \nabla^2 + \frac{\hbar^4}{8m^3 c^2} \nabla^4 + \frac{e^2}{8\pi\epsilon_0 r^3 m^2 c^2} \hat{\underline{S}} \cdot \hat{\underline{L}} - \left( \frac{\hbar^2}{4m^2 c^2} \underline{\underline{\nabla}} \nabla \right) \underline{\underline{\nabla}}$$

and

$$\hat{H} \psi = E \psi = (E - mc^2) \psi \quad - (41)$$

$$\quad \quad \quad - (42)$$

The wave equation is:

$$(\hat{H} - E) \underline{\underline{\nabla}} \psi = \underline{\underline{F}} \psi \quad - (43)$$

7)

The conventions used in this note are:

$$\mathbf{p} \rightarrow \mathbf{p} - e\mathbf{A}, \quad (44)$$

$$\mathbf{E} \rightarrow \mathbf{E} - e\phi, \quad (45)$$

$$\nabla = -e\phi \quad (46)$$

$$\phi = -\frac{e}{4\pi\epsilon_0 r} \quad (47)$$

$$\mathbf{E} = -\nabla \phi = -\frac{e}{4\pi\epsilon_0} \frac{\mathbf{r}}{r^3} \quad (48)$$

$$e\mathbf{E} = -\nabla V \quad (49)$$

So:

$$\begin{aligned} \hat{H}_{so} &= -\frac{e\hbar}{4m^2c^2} \underline{\sigma} \cdot \mathbf{E} \times \hat{\mathbf{p}} \\ &= \frac{e}{8\pi\epsilon_0 r^3 m^2 c^2} \underline{\hat{S}} \cdot \underline{\hat{L}} \quad (49) \end{aligned}$$

which is the same as eq. (9.3.6) of Atkins, 2nd. ed., with  $Z = 1$ .

Note that the convention used by Atkins is different because he defines:

$$V = -\frac{e^2}{4\pi\epsilon_0 r^2} = -e\phi, \quad (50)$$

$$\text{so } \phi = \frac{e}{4\pi\epsilon_0 r} \quad (\text{Atkins}) \quad (51)$$

This makes no difference to the final result (49).