

# THE SAGNAC EFFECT IN $m$ THEORY

by

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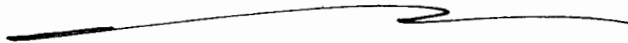
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## ABSTRACT

It is shown that the Sagnac effect is a method of experimental measurement of the  $m(r)$  function of  $m$  theory, and an experimental method of investigating the gravitational dependence of  $m(r)$ .

Keywords: ECE Unified field theory,  $m$  theory, Sagnac effect.

UFT 422



## 1. INTRODUCTION

In immediately preceding papers of this series {1 - 41} the m theory of natural philosophy has been developed from the infinitesimal line element of the most general spherically symmetric spacetime. In section 2 this line element is applied to the well known Sagnac interferometer, and it is shown that  $m(r)$  can be measured experimentally. Section 3 presents graphics for various  $m(r)$  functions. This short paper is based on Note 420(4) on [www.aias.us](http://www.aias.us).

## 2. DERIVATION OF THE SAGNAC EFFECT

Consider the infinitesimal line element of m theory in the plane polar coordinates

$$(r, \phi): \quad ds^2 = c^2 d\tau^2 = m(r)c^2 dt^2 - \frac{dr^2}{m(r)} - r^2 d\phi^2 \quad (1)$$

where  $\tau$  is the proper time and  $m(r)$  is any function of  $r$ . In the Einsteinian general relativity (EGR):

$$m(r) = 1 - \frac{2MG}{c^2 r} \quad (2)$$

where  $M$  is the attracting mass,  $G$  is the gravitational constant, and  $c$  is the speed of light in vacuo. From Einstein's equivalence principle an object such as a photon travelling at the speed of light is described by a null geodesic:

$$ds^2 = 0. \quad (3)$$

The distance  $r$  in the Sagnac effect is the radius of the rotating platform of the interferometer, so it does not change:

$$dr = 0. \quad (4)$$

It follows that:

$$m(r)c^2 dt^2 = r^2 d\phi^2. \quad (5)$$

The Sagnac effect is consequently:

$$d\phi \rightarrow d\phi + \Omega dt \quad (6)$$

where  $\Omega$  is the angular frequency of rotation of the platform. Eq. (6) is an example of the frame rotation theory developed in immediately preceding UFT papers. This concept is used for example in de Sitter precession and Thomas precession. It follows that:

$$m(r)^{1/2} dt = \frac{r}{c} (d\phi + \Omega dt). \quad (7)$$

Define the angular frequency of light traversing the Sagnac interferometer as:

$$\omega = \frac{c}{r} \quad (8)$$

and define the angular frequency of platform rotation by:

$$\Omega = \frac{v}{r}. \quad (9)$$

It follows that:

$$dt = \frac{\frac{1}{\omega} d\phi}{m(r)^{1/2} - \frac{\Omega}{\omega}} \quad (10)$$

So:

$$dt = \frac{d\phi}{m(r)^{1/2} \omega - \Omega}. \quad (11)$$

Integrating over the  $2\pi$  orbit of light traversing the perimeter of the platform:

$$t_1 = \frac{2\pi}{m(r)^{1/2} \omega - \Omega} \quad - (12)$$

For light propagating in the opposite sense:

$$t_2 = \frac{2\pi}{m(r)^{1/2} \omega + \Omega} \quad - (13)$$

It follows that:

$$\Delta t = t_1 - t_2 = \frac{4\pi \Omega}{m(r) \omega^2 - \Omega^2} \quad - (14)$$

Now use:

$$\omega \gg \Omega \quad - (15)$$

to find that:

$$\Delta t = \frac{4 A_r \Omega}{m(r) c^2} \quad - (16)$$

where the area of the platform is:

$$A_r = \pi r^2 \quad - (17)$$

Therefore  $m(\checkmark)$  can be measured experimentally in a high sensitivity Sagnac interferometer. The time difference is measured by interferometry, and the area of the platform is maximized by using many turns of a thin optical fibre.

This type of interferometer is compact and the time difference can be measured in the earth's gravitational field in a laboratory at sea level, and in a spacecraft under zero gravity conditions. This would measure the dependence of  $m(\checkmark)$  on gravity. Measurements of

the Sagnac time difference at different altitudes would reveal the dependence of  $m(r)$  on  $r$ .

EGR predicts that:

$$m(r) = 1 - \frac{2m_0}{c^2 r} \quad - (18)$$

where  $r$  is the radial coordinate. For a Sagnac interferometer on the earth's surface,  $r$  would be the radius of the earth according to EGR. The Schwarzschild radius of the earth is 0.09 metres so if  $r$  is interpreted as the radius of the platform:

$$m(r) \text{ (EGR)} = 1 - 0.09 / r \quad - (19)$$

and  $m(r)$  would depend on the radius  $r$  of a Sagnac platform. This can be easily investigated experimentally and is a test of EGR. The above derivation is also an experimental test of the Einstein equivalence principle in  $m$  space.

### 3. GRAPHICS OF THE SAGNAC EFFECT FOR VARIOUS M FUNCTIONS

Section by Dr. Horst Eckardt

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## 3 Graphics of the Sagnac effect for various m functions

According to Eq. (14) the time difference measured by a Sagnac interferometer rotated in two directions is

$$\Delta t = \frac{4\pi\Omega}{\omega^2 m(r) - \Omega^2} \quad (20)$$

where  $\omega$  is the angular frequency of the light and  $\Omega$  is that of mechanical rotation. The time difference will depend on the distance  $r$  from the gravitational centre if  $m(r)$  sufficiently differs from unity in the radial range investigated. Eq. (20) has been evaluated graphically for a demo system in Fig. 1. The model  $m$  function is that derived from Einsteinian theory:

$$m(r) = 1 - \frac{r_S}{r} \quad (21)$$

with so-called Schwarzschild radius

$$r_S = \frac{2MG}{c^2} \quad (22)$$

where  $M$  is the gravitating mass. From Fig. 1 it can be seen that  $\Delta t$  has a pole at  $r = r_S$ , however  $r_S$  normally lies inside the gravitating body so that we always have  $r \gg r_S$  where  $m(r)$  is nearly unity. Correspondingly, the time differences to be expected are small and are determined by the right hand side asymptotic value which is

$$\Delta t \rightarrow \frac{4\pi\Omega}{\omega^2 - \Omega^2}. \quad (23)$$

Values of  $\Delta t$  for some celestial bodies are listed in Table 1. We assumed  $\Omega = 10^4/\text{min}$  and  $\omega = 10^{15}/\text{s}$ , i.e. for a Sagnac interferometer with optical fibres.

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| body            | $m$ [kg]   | $r_S$ [m] | $r$ [m]    | $\Delta t$ [s] |
|-----------------|------------|-----------|------------|----------------|
| earth           | 5.97219e24 | 0.00887   | 6.371009e6 | 1.32e-26       |
| sun             | 1.98855e30 | 2953      | 6.95508e8  | 1.32e-26       |
| galactic centre | 8.36e36    | 1.24e10   | 1.24e11    | 1.46e-26       |

Table 1: Parameters of Sagnac effect for  $\omega = 10^{15}$ /s,  $\Omega = 10^4$ /min.

When measured at the surface of the earth and (hypothetically) at the surface of the sun, it is seen that the result is  $\Delta t = 1.32 \cdot 10^{-26}$  s in both cases. This means that the time difference is determined by the interferometric limit (23) and no dependence on  $r$  is detectable any more. Even when inspecting the case of the galactic centre, which is an extremely heavy star with a Schwarzschild radius of  $10^{10}$ m, the time difference would already be in saturation at the ten-fold distance of this radius.

In order to obtain well measurable time differences one would have to reduce the frequency of electromagnetic radiation in the interferometer drastically. The time difference depends on the inverse square of  $\omega$ . We have graphed the dependence of  $\Delta t$  from  $\omega$  for a fixed earth radius in Fig. 2. The curves are presented on logarithmic scales. Obviously the radiation frequency has to be lowered to the MHz range to obtain time differences in the range of  $10^{-8}$ s. For comparison we have added a curve for a different  $m$  function (exponential function) we used in preceding UFT papers:

$$m(r) = 2 - \exp\left(\log(2) \exp\left(-\frac{r}{R}\right)\right). \quad (24)$$

We had to increase the parameter  $R$  to  $10^7$ m to obtain a visible difference in the diagram. Practical measurements of  $m(r)$  seem to be a hard challenge for an experiment.

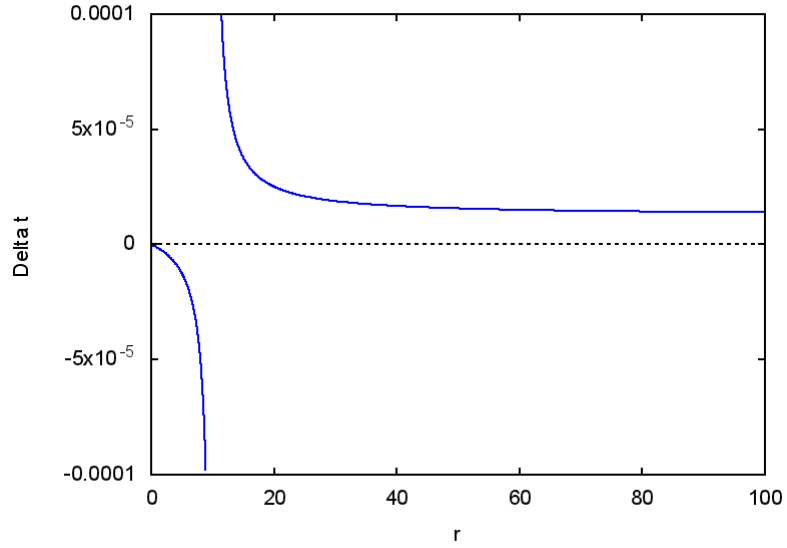


Figure 1: Principal dependence of Sagnac effect on  $r$ .

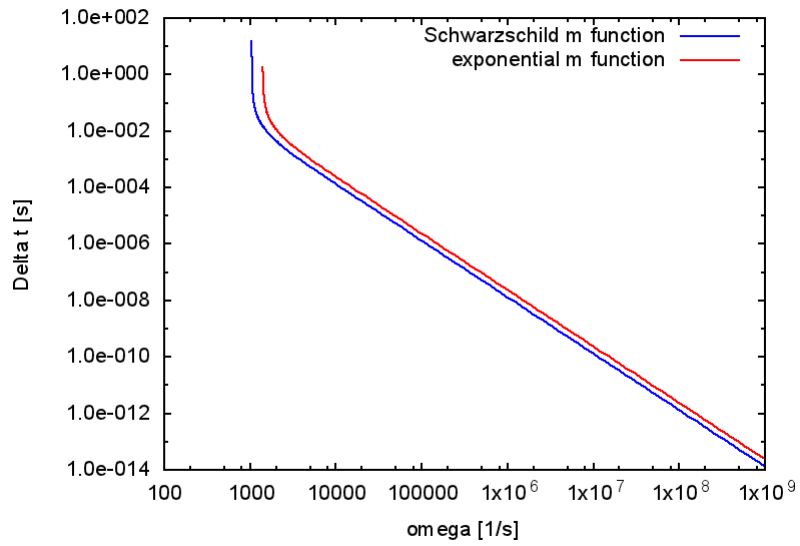


Figure 2: Dependence of Sagnac effect on  $\omega$  at earth surface.



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