

1) Revised Method for Paper 85

The radiative correction to the Schrödinger equation of the H atom is, to first order in d :

$$-\frac{\hbar^2}{2m} \nabla^2 \left(1 + \frac{d}{2\pi}\right) \psi - \frac{e^2}{4\pi\epsilon_0 r} \psi = E \psi. \quad (1)$$

This is equivalent to:

$$-\frac{\hbar^2}{2m} \nabla^2 \psi - \frac{e^2}{4\pi\epsilon_0 (r + r(\text{vac}))} \psi = E \psi. \quad (2)$$

Assume that $r(\text{vac})$ is small, so:

$$\psi = \psi_0 \quad (3)$$

where ψ_0 is the H atom wave function in the absence of radiative corrections. Now subtract eq. (2) from eq. (1):

$$-\frac{\hbar^2 d}{4\pi m} \nabla^2 \psi_0 = \frac{e^2}{4\pi\epsilon_0} \left(\frac{1}{r} - \frac{1}{r + r(\text{vac})} \right) \psi_0 \quad (4)$$

i.e.

$$\nabla^2 \psi_0 = -\frac{4\pi m e^2}{\hbar^2} \left(\frac{1}{r} - \frac{1}{r + r(\text{vac})} \right) \psi_0 \quad (5)$$

Evaluating this equation by computer algebra gives:

$$2) \quad \frac{1}{r + r_{2p}(\text{vac})} - \frac{1}{r + r_{2s}(\text{vac})} = \frac{1}{4\pi^{3/2}} \frac{\hbar}{mc} \frac{1}{r^2} \quad - (6)$$

The change in potential energy is:

$$\Delta V = d \left(\frac{1}{r + r_{2p}(\text{vac})} - \frac{1}{r + r_{2s}(\text{vac})} \right) \text{ cm}^{-1} \quad - (7)$$

$$= \frac{d}{4\pi^{3/2}} \frac{\hbar}{mc} \frac{1}{r^2} \quad - (8)$$

The change in total energy is:

$$\Delta E = \frac{r}{2n^2 a} \Delta V = \frac{d}{32\pi^{3/2}} \frac{\hbar}{mc} \frac{1}{a} \cdot \frac{1}{r}$$

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$$\Delta E = \left(\frac{1}{32\pi^{3/2}} \frac{d}{a} \frac{\hbar}{mc} \right) \frac{1}{r} \text{ cm}^{-1} \quad - (9)$$

$$= 0.0353 \text{ cm}^{-1}$$

This gives: $r = 8.46 \times 10^{-4} \text{ m} \quad - (10)$

From eq. (6):

$$\frac{r_{2s}(\text{vac}) - r_{2p}(\text{vac})}{(r + r_{2p}(\text{vac}))(r + r_{2s}(\text{vac}))} = \frac{1}{4\pi^{3/2}} \frac{\hbar}{mc} \frac{1}{r^2} \quad - (11)$$

Eq. (3) is consistent with the assumption:

3)

$$r > r_{2s} \sim r_{2p}, \quad - (12)$$

$$r \gg (r_{2s}(\text{vac}) - r_{2p}(\text{vac})) \quad - (13)$$

So from eq. (11):

$$r_{2s}(\text{vac}) - r_{2p}(\text{vac}) \sim \frac{1}{4\pi^{3/2}} \frac{\hbar}{mc} \quad - (14)$$

i.e. $r_{2s}(\text{vac}) - r_{2p}(\text{vac}) \sim \frac{1}{8\pi^{5/2}} \frac{\hbar}{mc} \quad - (15)$

where $\lambda_c = \text{Compton wavelength} = \frac{\hbar}{mc} = 2.426 \times 10^{-12} \text{ m}$

so $r_{2s}(\text{vac}) - r_{2p}(\text{vac}) = 1.74 \times 10^{-14} \text{ m} \quad - (16)$

and indeed this is much less than r (eq. (10))

This result can be compared with the classical radius of the electron:

$$r(\text{classical}) = \frac{1}{4\pi\epsilon_0} \frac{e^2}{mc^2} = 2.818 \times 10^{-15} \text{ m} \quad - (17)$$

and the Bohr radius:

$$a = 5.292 \times 10^{-11} \text{ m} \quad - (18)$$

From these results it is seen that the vacuum perturbation of the electron is expected, i.e. much less than the Bohr radius and a little greater than its classical radius