

# Some Background Notes for Paper 61

The Poisson equation for the radial component of  $\epsilon_0$  system  $(r, \theta, \phi)$  is:

$$\frac{d^2 \phi}{dr^2} + \frac{2}{r} \frac{d\phi}{dr} = -\frac{\rho}{\epsilon_0} \quad - (1)$$

In ECE theory, this becomes:

$$\frac{d^2 \phi}{dr^2} + \frac{2}{r} \frac{d\phi}{dr} = -\frac{\rho}{\epsilon_0} + A \left( \frac{1}{r} \frac{d\phi}{dr} + \frac{\phi}{r^2} \right)$$
$$= \frac{1}{\epsilon_0} (-\rho + \rho_{\text{eff}}) \quad - (2)$$

where:

$$\rho_{\text{eff}} = A \epsilon_0 \left( \frac{1}{r} \frac{d\phi}{dr} + \frac{\phi}{r^2} \right) \quad - (3)$$

The A term decreases  $-\rho$ , i.e. the negative potential energy of the Coulomb law is added by a positive potential energy inputted from spacetime:

$$V_{\text{eff}} = e \phi_{\text{eff}} \quad - (4)$$

where  $\phi_{\text{eff}}$  is the potential caused by  $\rho_{\text{eff}}$ .  
The net effect is that the electron is attracted less strongly to the proton than in the absence of A.

A. When:

$$\rho_{\text{eff}} > \rho \quad - (5)$$

2) If repulsion exceeds the attraction and the H atom is ionized. The ionization energy of H is 13.6 eV, and when  $V_{eff}$  is greater than this amount, the electron breaks free of the proton.

A particular solution of eq. (2) is given by:

$$A = 2 \quad - (6)$$

when

$$\frac{d^2 \phi}{dr^2} = \underbrace{-\frac{e}{\epsilon_0}}_{\text{attraction}} + \underbrace{\frac{2\phi}{r^2}}_{\text{repulsion}} \quad - (7)$$

This is an undamped resonant oscillator if  $\phi$  is initially oscillatory. If  $\phi$  is made very large and positive, and if  $e$  is positive, the potential energy in eq. (4) is positive.

If the second term on the right hand side of eq. (7) dominates, then:

$$\phi \sim \beta / r \quad - (8)$$

is a possible solution, where  $\beta$  must have the S.I. units of  $e / (4\pi\epsilon_0)$ . In this simple approximation, the potential energy of the H atom is simply made more positive:

$$V = -\frac{e^2}{4\pi\epsilon_0 r} + \frac{e\beta}{r} \quad - (9)$$

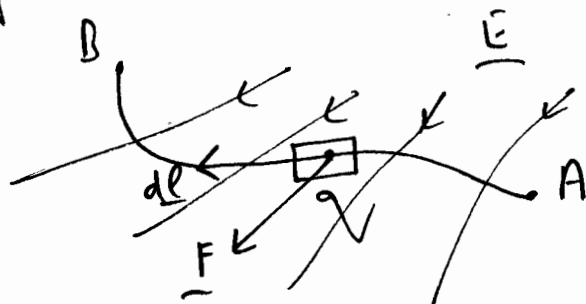
3) and if  $\beta$  is large enough there will be no negative attraction energy between proton and electron.

### Sign Conventions of the Standard Model

The scalar potential in electrostatics is defined by:

$$\underline{E} = -\underline{\nabla} \phi \quad (10)$$

in the standard model. The work done on a test charge  $q$  in moving it from point A to point B in the presence of  $\underline{E}$  is sketched as follows:



The force acting on a charge is:

$$\underline{F} = q \underline{E} \quad (11)$$

The work done is:

$$W = - \int_A^B \underline{F} \cdot d\underline{l} = -q \int_A^B \underline{E} \cdot d\underline{l} \quad (12)$$

The minus sign appears because the work done on the charge is being calculated. This work done is calculated against the action of the electric field.

With the definition (10):

$$\begin{aligned}
 W &= q \int_A^B \underline{\nabla} \phi \cdot d\underline{\ell} \\
 &= q \int_A^B d\phi \\
 &= q (\phi_B - \phi_A) \quad \text{--- (13)}
 \end{aligned}$$

This is the potential energy difference between B and A.

The line integral of the electric field between two points is independent of the path and is the negative of the potential difference:

$$\int_A^B \underline{E} \cdot d\underline{\ell} = -(\phi_B - \phi_A) \quad \text{--- (14)}$$

If the path is closed there is no potential energy difference, so:

$$\oint \underline{E} \cdot d\underline{\ell} = 0. \quad \text{--- (15)}$$

Stokes' theorem states that:

$$\oint_C \underline{A} \cdot d\underline{\ell} = \int_S \underline{\nabla} \times \underline{A} \cdot \underline{n} \, da \quad \text{--- (16)}$$

$$\text{so:} \quad \underline{\nabla} \times \underline{E} = \underline{0}. \quad \text{--- (17)}$$

If we go through these well known formulae in ECE theory, the electric field

5) is changed to:

$$\underline{E} = -\underline{\nabla} \phi + \phi \underline{\omega} \quad - (18)$$

and it becomes less negative by an amount:

$$\boxed{\underline{E}_{\omega} = \phi \underline{\omega}}$$

-(19)

(volt m<sup>-1</sup>)

So the work done in eq. (12) is changed by

an amount:

$$\underline{W}_{\omega} = -q \int_A^B \underline{E}_{\omega} \cdot d\underline{l} \quad - (20)$$

$$= -q \int_A^B \phi \underline{\omega} \cdot d\underline{l}$$

So eq. (13) is changed to:

$$\underline{W} = q(\phi_B - \phi_A) - \underline{W}_{\omega} \quad - (21)$$

This means that less work is needed to move  $q$  from  $A$  to  $B$  against the electric field.

The electric field is weaker and the Coulombic attraction <sup>or repulsion</sup> is weaker. This has

assumed that  $\underline{\omega}$  is positive.

Repulsion between two like charges produce a positive  $\phi_B - \phi_A$ , attraction between unlike charges a negative  $\phi_B - \phi_A$

For the Coulomb law of the standard model

6) attraction between a proton of charge  $+e$  and an electron of charge  $-e$  is described by the negative potential difference:

$$V = -\frac{e^2}{4\pi\epsilon_0 r} \quad - (22)$$

A positive  $\underline{\omega}$  makes this less negative:

$$\underline{F} = -e \underline{E}, \quad - (23)$$

(corresponding to eq. (22), is changed to:

$$\underline{F} = -e (\underline{E} - \phi \underline{\omega}), \quad - (24)$$

i.e

$$\underline{F} = -e (\underline{E} - \underline{E}_\omega) \quad - (24a)$$

The repulsive potential energy difference:

$$V_\omega = +q \int_A^B \phi \underline{\omega} \cdot d\underline{l} \quad - (25)$$

is inputted for spacetime if  $\underline{\omega}$  is positive.

Negative  $V$  denotes attraction,  
positive  $V$  denotes repulsion.