

172(5): Simple Solution of the ECE Fermion Equation

Consider for simplicity motion in the  $z$  axis, then the ECE fermion equation is:

$$\sigma^0 \hat{E} \psi - c \hat{p}_z \sigma^3 \psi \sigma^3 = \sigma^1 mc^2 \psi \quad (1)$$

in which:  $\hat{E} = i\hbar \frac{\partial}{\partial t}$ ,  $\hat{p}_z = -i\hbar \frac{\partial}{\partial z} \quad (2)$

and  $\sigma^0 = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$ ,  $\sigma^3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$ ,  $\sigma^1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad (3)$

The eigenfunction  $\psi$  is a bispinor:

$$\psi = \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} \quad (4)$$

Eq. (1) may be solved analytically as follows. First note that:

$$\sigma^1 \psi = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} = \begin{bmatrix} \psi_1^L & \psi_2^L \\ \psi_1^R & \psi_2^R \end{bmatrix} \quad (5)$$

and

$$\sigma^3 \psi \sigma^3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \quad (6)$$

$$= \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} \psi_1^R & -\psi_2^R \\ \psi_1^L & -\psi_2^L \end{bmatrix} = \begin{bmatrix} \psi_1^R & -\psi_2^R \\ -\psi_1^L & \psi_2^L \end{bmatrix} \quad (7)$$

so eq. (1) is:

$$\hat{E} \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} - c \hat{p}_z \begin{bmatrix} \psi_1^R & -\psi_2^R \\ -\psi_1^L & \psi_2^L \end{bmatrix} = mc^2 \begin{bmatrix} \psi_1^L & \psi_2^L \\ \psi_1^R & \psi_2^R \end{bmatrix} \quad (8)$$

i.e.

$$\begin{aligned} \hat{E} \psi_1^R - c \hat{p}_z \psi_1^R &= mc^2 \psi_1^L & (9) \\ \hat{E} \psi_2^R + c \hat{p}_z \psi_2^R &= mc^2 \psi_2^L & (10) \\ \hat{E} \psi_1^L + c \hat{p}_z \psi_1^L &= mc^2 \psi_1^R & (11) \\ \hat{E} \psi_2^L - c \hat{p}_z \psi_2^L &= mc^2 \psi_2^R & (12) \end{aligned}$$

The Dirac equation corresponding to eq. (1) is:

$$(\gamma^0 \hat{E} + c \gamma^3 \hat{p}_3) \psi_0 = mc^2 \psi_0 \quad - (13)$$

where:  $\gamma^0 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$ ,  $\gamma^3 = \begin{bmatrix} 0 & -\sigma^3 \\ \sigma^3 & 0 \end{bmatrix} \quad - (14)$

and  $p_3 = -p_2 \quad - (15)$

Here  $\psi_0$  is the Dirac spinor. Written out in full, eq. (13) is:

$$\begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \hat{E} \begin{bmatrix} \psi_1^R \\ \psi_2^R \\ \psi_1^L \\ \psi_2^L \end{bmatrix} - c \begin{bmatrix} 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{bmatrix} \hat{p}_2 \begin{bmatrix} \psi_1^R \\ \psi_2^R \\ \psi_1^L \\ \psi_2^L \end{bmatrix} = mc^2 \begin{bmatrix} \psi_1^R \\ \psi_2^R \\ \psi_1^L \\ \psi_2^L \end{bmatrix} \quad - (16)$$

which gives eqs. (9) to (12) again.

The advantage of eq. (1) is that it can be written in the same overall structure for the general case, and uses  $2 \times 2$  matrices only. One of the advantages of the Dirac eq. (13) was  $4 \times 4$  matrices which are not entirely desirable. Eq. (13) is the chiral representation. There is also a standard representation.

The structure of eq. (1) is:

$$\boxed{\sigma^0 \hat{p}_0 \psi + \sigma^3 \hat{p}_3 \psi = \sigma^1 mc \psi} \quad - (17)$$

where  $\hat{p}_0 = \hat{E} / c \quad - (18)$

and is a type of tensor product of the LHS of Pauli matrices, which are basis elements. Eq. (17)

Therefore use only the well defined Pauli matrices:  
 $\sigma^0 = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$ ,  $\sigma^1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$ ,  $\sigma^2 = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$ ,  $\sigma^3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$  - (19)  
 and in eq. (17) the eigenvalue  $mc$  is always positive.  
 Eq. (17) can be developed greatly for quantum field  
 theory, particle theory and quantum electrodynamics.

By inspection, a simple solution of the set  
 of equations (9) to (12) can be found. First consider  
 equations (9) and (11) in position representation:

$$\left( \frac{\partial}{\partial t} + c \frac{\partial}{\partial z} \right) \psi_1^R = -i\omega_0 \psi_1^R \quad - (19)$$

$$\left( \frac{\partial}{\partial t} - c \frac{\partial}{\partial z} \right) \psi_1^L = -i\omega_0 \psi_1^L \quad - (20)$$

$$\omega_0 = mc^2 / \hbar \quad - (21)$$

where  
 is the rest frequency of the fermion.

$$\text{If: } \psi_1^R = \exp(-i(\omega_0 t + \kappa z)) \quad - (22)$$

$$\psi_1^R = \frac{\omega_0}{\omega_0 + c\kappa} \psi_1^L \quad - (23)$$

$$\psi_1^L = \frac{\omega_0 + c\kappa}{\omega_0} \psi_1^R \quad - (24)$$

from eq. (19).

$$\text{If } \psi_1^L = \exp(i(\omega_0 t + \kappa z)) \quad - (25)$$

i) then

$$\psi_1^L = \frac{\omega_0}{c\kappa - \omega_0} \psi_1^R \quad - (26)$$

from eq. (20). From eqs. (24) and (26):

$$\frac{\omega_0 + c\kappa}{\omega_0} = \frac{\omega_0}{c\kappa - \omega_0} \quad - (27)$$

so  $\omega_0^2 = (c\kappa + \omega_0)(c\kappa - \omega_0) \quad - (28)$

i. e.  $2\omega_0^2 = c^2 \kappa^2 \quad - (29)$

$$\omega_0 = \left( \frac{c}{\sqrt{2}} \right) \kappa \quad - (30)$$

This is a possible solution with:

$$\begin{aligned} \psi_1^R &= e^{-i\phi} \\ \psi_1^L &= e^{i\phi} \\ \phi &= \omega_0 t + \kappa z \end{aligned} \quad - (31)$$

Another solution is:

$$\begin{aligned} \psi_1^R &= \exp(-i(\omega_0 t - \kappa z)) \\ \psi_1^L &= \exp(i(\omega_0 t - \kappa z)) \end{aligned} \quad - (32)$$

which again gives eq. (30).

For these solutions the real parts of  $\psi_1^R$  and  $\psi_1^L$  are the same, and from eqs. (9) to (12) it is seen that the energy eigenvalue  $mc^2$  is change positive.

## Anti-Fermion Equations

This is found by the parity operations:

$$\hat{P}(p_z) = -p_z; \quad \hat{P}(\psi_1^R) = \psi_1^L; \quad \hat{P}(\psi_2^R) = \psi_2^L;$$

$$\hat{P}(\psi_1^L) = \psi_1^R; \quad \hat{P}(\psi_2^L) = \psi_2^R. \quad \text{So it is found that:}$$

$$\hat{P} \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix} = \begin{bmatrix} \psi_1^L & \psi_2^L \\ \psi_1^R & \psi_2^R \end{bmatrix} \quad - (33)$$

i.e.

$$\boxed{\hat{P} \psi = \sigma^1 \psi} \quad - (34)$$

so the ECE fermion equation (17) is:

$$\boxed{\sigma^0 \hat{P}_0 \psi \sigma^0 + \sigma^3 \hat{P}_3 \psi \sigma^3 = mc \hat{P} \psi} \quad - (35)$$

The ECE anti-fermion equation is: - (36)

$$\hat{E} \begin{bmatrix} \psi_1^L & \psi_2^L \\ \psi_1^R & \psi_2^R \end{bmatrix} + c \hat{P}_2 \begin{bmatrix} \psi_1^L & -\psi_2^L \\ -\psi_1^R & \psi_2^R \end{bmatrix} = mc^2 \begin{bmatrix} \psi_1^R & \psi_2^R \\ \psi_1^L & \psi_2^L \end{bmatrix}$$

i.e.

$$\boxed{\sigma^0 \hat{P}_0 (\hat{P} \psi) \sigma^0 - \sigma^3 \hat{P}_3 (\hat{P} \psi) \sigma^3 = mc \psi} \quad - (37)$$

The energy eigenvalue  $mc^2$  is again positive.