

100(14) : Potential Equivalent of Einstein Postulate

The original Einstein postulate of 1915 is well known to be:

$$D^\mu G_{\mu\nu} = k D^\mu T_{\mu\nu} = 0, \quad - (1)$$

where:

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}. \quad - (2)$$

Here $G_{\mu\nu}$ is the Einstein tensor, $R_{\mu\nu}$ is the Ricci tensor, R is the scalar curvature, $g_{\mu\nu}$ is the symmetric metric, D^μ indicates covariant derivative, k is Einstein's constant and $T_{\mu\nu}$ is the symmetric canonical energy-momentum density tensor. In eq. (1) the second Bianchi identity is:

$$D^\mu G_{\mu\nu} = 0 \quad - (3)$$

and Noether's Theorem is:

$$D^\mu T_{\mu\nu} = 0. \quad - (4)$$

From eq. (1) Einstein used the particular solution:

$$G_{\mu\nu} = k T_{\mu\nu}. \quad - (5)$$

In paper 88 it was shown that eq. (3) is true if and only if the tensor is identically zero. In the presence of torsion:

$$D_\mu T^{\kappa\mu\nu} = -R^{\kappa\mu\nu} \quad - (6)$$

$$D_\mu \tilde{T}^{\kappa\mu\nu} = -\tilde{R}^{\kappa\mu\nu} \quad - (7)$$

2) Eqs. (6) and (7) are particular solutions of the true Bianchi identity:

$$D \wedge T^a := R^a{}_b \wedge q^b \quad - (8)$$

In the notation of paper 88 it was shown that eq. (8) is the only Bianchi identity. The usually named second Bianchi identity is a special case of:

$$D \wedge (D \wedge T^a) := D \wedge (R^a{}_b \wedge q^b) \quad - (9)$$

which may be written as:

$$(D \wedge R)_{cyclic} := (D(D \wedge T))_{cyclic} \quad - (10)$$

The special case used by Einstein is:

$$D \wedge (R^a{}_b \wedge q^b) = 0 \quad - (11)$$

i.e. $D \wedge R^a{}_b = 0 \quad - (12)$

which is also the special case considered by (small).

In tensor notation eq. (8) is:

$$D_{\mu} \tilde{T}^{\mu\nu} = -\tilde{R}^{\mu\nu} \quad - (13)$$

A well-defined Hodge dual of eq. (8) is:

$$D \wedge \tilde{T}^a := \tilde{R}^a{}_b \wedge q^b \quad - (14)$$

which in tensor notation is:

$$D_{\mu} T^{\kappa\mu\nu} = -R^{\kappa\mu\nu}{}_{\mu} \quad - (15)$$

It follows from eq. (15) that if $R^{\kappa\mu\nu}$ is non-zero then the torsion tensor $T^{\kappa\mu\nu}$ is also non-zero.
 The Ricci cyclic equation is:

$$R^a{}_b \wedge \eta^b = 0 \quad - (16)$$

which implies: $D \wedge T^a = 0. \quad - (17)$

Eq (17) does not imply that:

$$T^a = 0, \quad - (18)$$

and as was shown in paper 93, $R^{\kappa\mu\nu}$ is zero only for Ricci flat spacetimes, i.e. only for:

$$G_{\mu\nu} = 0, \quad T_{\mu\nu} = 0. \quad - (19)$$

One can only say that for metrics that obey the Ricci cyclic equation:

$$D_{\mu} \tilde{T}^{\mu\nu} = -\tilde{R}^{\kappa\mu\nu}{}_{\mu} = 0, \quad - (20)$$

but $D_{\mu} T^{\kappa\mu\nu} = -R^{\kappa\mu\nu}{}_{\mu} \neq 0, \quad - (21)$

in general. However, for the Christoffel connection

$$4) \quad T^{\mu\nu} = \Gamma^{\mu\nu} - \Gamma^{\nu\mu} = 0. \quad - (22)$$

This is the special case considered by Einstein and Hilbert. In this special case it is possible to state only that if

$$\tilde{R}^{\mu\nu} = 0 \quad - (22)$$

then

$$R^{\mu\nu} \neq 0 \quad - (23)$$

in general. For a Christoffel connection:

$$T^{\mu\nu} = 0 \quad - (24)$$

$$\tilde{T}^{\mu\nu} = 0. \quad - (25)$$

Eq. (22) is equivalent to:

$$R^{\mu\nu} + R^{\nu\mu} = 0 \quad - (26)$$

and Eq. (23) is:

$$\tilde{R}^{\mu\nu} + \tilde{R}^{\nu\mu} \neq 0 \quad - (27)$$

in general.

So the well known EH equation (5) is a very special case that loses information.

5) If a Christoffel connection is used, eq. (6) can only be rigorously true in a Ricci flat spacetime, where:

$$D_{\mu} T^{\kappa\mu\nu} = -R^{\kappa\mu\nu}{}_{\mu} = 0. \quad (28)$$

In general, eq. (28) indicates that the Christoffel connection cannot be used in general, because for a Christoffel connection:

$$T^{\kappa\mu\nu} = 0 \text{ but } R^{\kappa\mu\nu}{}_{\mu} \neq 0 \quad (29)$$

resulting in a contradiction.

This contradiction is present in the EH

theory of gravitation.

In paper 93, eq. (6) was interpreted as the balance of different fields, electromagnetic and gravitational:

$$A^{(e)} (D_{\mu} T^{\kappa\mu\nu})^{\text{electromagnetic}} = -A^{(g)} (R^{\kappa\mu\nu}{}_{\mu})^{\text{gravitational}} \quad (30)$$

so $T^{\kappa\mu\nu}$ and $R^{\kappa\mu\nu}{}_{\mu}$ refer to two different fields. This is an approximation. In paper 93

the Christoffel connection was used to compute

$$R^{\kappa\mu\nu}{}_{\mu} \neq 0 \quad (31)$$

6) in general. This means that there is no gravitation tensor present:

$$(\mathbb{T}^{\mu\nu})_{\text{gravitation}} = 0 \quad - (32)$$

but there is an electromagnetic tensor present.

If we are considering one field only, then eq. (28) prohibits the use of Φ (Christoffel symbol). To describe gravitation correctly, the more general gamma connection must be used, i.e.

$$\boxed{\Gamma^{\lambda}_{\mu\nu} \neq \Gamma^{\lambda}_{\nu\mu}} \quad - (33)$$

It is reasonable then to assume that the more general conservation law is:

$$\boxed{D_{\mu} \mathbb{T}^{\mu\nu} = k D_{\mu} \mathbb{J}^{\mu\nu} = -R^{\lambda}_{\mu}} \quad - (34)$$

A particular solution is:

$$\boxed{\mathbb{T}^{\mu\nu} = k \mathbb{J}^{\mu\nu}} \quad - (35)$$

where $\mathbb{J}^{\mu\nu}$ is the canonical angular energy-momentum density tensor:

$$\mathbb{J}^{\mu\nu} = -\frac{1}{2} (\mathbb{T}^{\mu\alpha} x^{\nu} - \mathbb{T}^{\nu\alpha} x^{\mu}) \quad - (36)$$