

8(5): Precession from Various Metrics

In this note we consider the rotation of various metrics in order to deduce the effect of the precession of these metrics on the Thomas precession and Thomas half. First consider the Minkowski metric in polar coordinates:

$$ds^2 = c^2 dt^2 - dr^2 - r^2 d\phi^2 \quad (1)$$

spherical polar coordinates:

$$ds^2 = c^2 dt^2 - dr^2 - r^2 d\theta^2 - r^2 \sin^2 \theta d\phi^2 \quad (2)$$

the rotation is defined by:

$$d\phi' = d\phi + \omega dt \quad (3)$$

both cases.

From eq. (3) in eq. (1):

$$ds^2 = (c^2 - r^2 \omega^2) dt^2 - 2\omega r^2 d\phi dt - (dr^2 + r^2 d\phi'^2) \quad (4)$$

the velocity of rotation is defined by

$$v = \omega r \quad (5)$$

$$\text{so } ds^2 = \left(1 - \frac{v^2}{c^2}\right) (c^2 dt^2 - 2r^2 \Omega d\phi dt) - (dr^2 + r^2 d\phi'^2) \quad (6)$$

where

$$\Omega = \omega \left(1 - \frac{v^2}{c^2}\right)^{-1} \quad (7)$$

known as the Thomas angle (UFT 110). The definition of

$$\Omega \text{ follows from } \omega = \frac{v}{r} \rightarrow \frac{v'}{r'} = \frac{v}{r/\gamma} = \gamma^2 \frac{v}{r} \quad (8)$$

where

$$\gamma^2 = \left(1 - \frac{v^2}{c^2}\right)^{-1} \quad (9)$$

so

$$\omega \rightarrow \Omega = \frac{\omega}{\left(1 - \frac{v^2}{c^2}\right)} \quad (10)$$

A.K.D.

It follows that:

$$dt' = \left(1 - \frac{v^2}{c^2}\right)^{1/2} dt \quad (11)$$

and that the Thomas precession is:

$$\Delta\phi_g = \Omega dt' - \omega dt \quad (12)$$

$$= 2\pi(\gamma - 1)$$

In spherical polar coordinates, eqs. (2) and (3) give:

$$ds^2 = (c^2 - r^2 \sin^2 \theta \omega^2) dt^2 - 2r^2 \sin^2 \theta \omega d\phi dt - dr^2 - r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \quad (13)$$

In spherical polar coordinates:

$$v^2 = r^2 \omega^2 \sin^2 \theta \quad (14)$$

(15)

so the overall effect is:

$$c^2 dt^2 \rightarrow \left(1 - \frac{v^2}{c^2}\right) (c^2 dt^2 - 2r^2 \sin^2 \theta \Omega d\phi dt)$$

In plane polar coordinates the overall effect is:

$$c^2 dt^2 \rightarrow \left(1 - \frac{v^2}{c^2}\right) (c^2 dt^2 - 2\Omega r^2 d\phi dt) \quad (16)$$

In L.O. coordinate systems the Thomas precession is given by eq. (12).

The Thomas precession is true in any coordinate system.

Now consider the above calculations for the general, spherical, metric in plane polar coordinates:

$$ds^2 = m(r, t) c^2 dt^2 - \frac{dr^2}{m(r, t)} - r^2 d\phi^2 \quad (17)$$

and in spherical polar:

$$ds^2 = m(r, t) c^2 dt^2 - \frac{dr^2}{m(r, t)} - r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \quad (18)$$

the effect of the rotation (3) on metrics (17) and (18) is to modify the Thomas precession to:

$$\Delta\phi = 2\pi \left[\left(m(r, t) - \frac{v^2}{c^2} \right)^{-1/2} - 1 \right] \quad (19)$$

provided that:

$$\Omega := \omega \left(m(r, t) - \frac{v^2}{c^2} \right)^{-1} \quad (20)$$

The function $m(r, t)$ for various type of metric is given in Table 1.

Metric	$m(r, t)$
1) Schwarzschild	$1 - \frac{2mG}{c^2 r}$
2) Kerr Newman and Reissner-Nordström	$1 - \frac{2mG}{c^2 r} + \frac{Q^2}{r^2}$
3) Einstein Rosen Bridge	$1 - \frac{2mG}{c^2 r} - \frac{e^2}{r^2}$
4) Static de Sitter	$1 - \frac{r}{\Lambda^2}$
5) Reissner Weyl	$1 - \frac{A}{r} - \frac{\kappa B}{2r^2}$

Many of these metrics are given in chapter four of "FT301, Criticism of the Einstein Field Equation", chapter four, and shown to be incorrect due to neglect of torsion. The effect of metrics (1) to (5) on the velocity v^2 of the Thomas precession give the effect of gravity in the solid state model.

1) Schwarzschild Metric

$$v^2 \rightarrow v^2 + \frac{2mG}{r} \quad - (21)$$

2) Kerr Newmann

$$\frac{v^2}{c^2} \rightarrow \frac{v^2}{c^2} + \frac{2mG}{rc^2} - \frac{Q^2}{r^2} \quad - (22)$$

3) Einstein-Rosen Bridge

$$\frac{v^2}{c^2} \rightarrow \frac{v^2}{c^2} + \frac{2mG}{rc^2} + \frac{E^2}{r^2} \quad - (23)$$

4) Static de Sitter

$$\frac{v^2}{c^2} \rightarrow \frac{v^2}{c^2} + \frac{r^2}{d^2} \quad - (24)$$

5) Reissner-Weyl

$$\frac{v^2}{c^2} \rightarrow \frac{v^2}{c^2} + \frac{A}{r} + \frac{\kappa B^2}{2r^2} \quad - (25)$$

In none of the cases (21) to (25) are the predictions of the standard model verified. This becomes clear as soon as it is realized that the Thomas half plays a fundamental role in physics.

FLRW Metric

- (26)

This is described by:

$$ds^2 = c^2 dt^2 - a^2(t) \left(\frac{dr^2}{1 - kr^2} \right) - r^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

and for $k=0$ gives a flat universe for $k=-1$ an open universe and for $k=1$ a closed universe. Eqs. (3) and (26) give the effect of these universes on the Thomas half, and again, no such effect is observed.