

# 399(1) : Spin Correction Response due to Vacuum Fluctuations

Assume that :

$$\underline{\delta r} = \underline{\delta r}(0) \cdot \exp(-i \omega_0 t) \quad - (1)$$

where  $\underline{\delta r}$  is the vacuum fluctuation, taking place with angular frequency  $\omega$  and wave vector  $\underline{r}$ . The vacuum force is then :

$$\underline{F}(r) = -m \omega_0^2 \underline{\delta r} = m \frac{d^2 \underline{\delta r}}{dt^2} \quad - (2)$$

where  $\underline{E}(\text{vac})$  is the fluctuating vacuum field.

Therefore :

$$\underline{E}(\text{vac}) = \frac{m}{e} \omega_0^2 \underline{\delta r}(0) e^{-i \omega_0 t} \quad - (3)$$

However, in ECE2 :

$$\underline{E}(\text{vac}) = \underline{\omega} \phi_0 \quad - (4)$$

where  $\underline{\omega}$  is the spin correction vector and  $\phi_0$  is the potential in the absence of the vacuum. So :

$$\underline{\omega} \phi_0 = \frac{m}{e} \omega_0^2 \underline{\delta r}(0) e^{-i \omega_0 t} \quad - (5)$$

and

$$\frac{d^2}{dt^2} (\underline{\omega} \phi_0) = \frac{m}{e} \omega_0^4 \underline{\delta r}(0) e^{-i \omega_0 t}$$

$$= A e^{-i \omega_0 t} \quad - (6)$$

where

$$\underline{A} = \frac{m}{e} \omega_0^4 \underline{\delta r}(0) = \text{constant} \quad - (7)$$

2) Therefore:

$$\frac{\omega}{\omega} \frac{d^2 \phi_0}{dt^2} + 2 \frac{d\phi_0}{dt} \frac{d\omega}{dt} + \left( \frac{d^2 \omega}{dt^2} \right) \phi_0 = A e^{-i\omega t} \quad (8)$$

For  $\phi_0$  X component:

$$\frac{d^2 \phi_0}{dt^2} + \frac{2}{\omega_x} \frac{d\omega_x}{dt} \frac{d\phi_0}{dt} + \frac{1}{\omega_x} \frac{d^2 \omega_x}{dt^2} \phi_0 = \frac{A e^{-i\omega t}}{\omega_x} \quad (9)$$

and similarly for Y and Z.

For Euler Bernoulli resistance to occur eq.

(9) has to be of  $\phi_0$  format:

$$\ddot{x} + 2\beta \dot{x} + \omega_1^2 x = A \cos \omega t \quad (10)$$

so

$$\frac{2}{\omega_x} \frac{d\omega_x}{dt} = 2\beta \quad (11)$$

$$\omega_1^2 = \frac{1}{\omega_x} \frac{d^2 \omega_x}{dt^2} \quad (12)$$

Therefore  $\frac{d^2 \omega_x}{dt^2} = \omega_1^2 \omega_x \quad (13)$

and  $\frac{d\omega_x}{dt} = \beta \omega_x \quad (14)$

A solution of eq. (13) is:

$$\omega_x = \omega(0) \exp(i(\omega t - \underline{v} \cdot \underline{r})) \quad (15)$$

Using this solution,  $\beta$  is imaginary, so  $\beta$  is zero

in the physical interpretation, using the usual rule of physical part is the real part.

So eq. (9) becomes: -(16)

$$\frac{d^2 \phi_0}{dt^2} + \omega_1^2 \phi_0 = A \cos(\omega_0 t + \omega_1 t - \underline{\kappa \cdot r})$$

Taking the real part:

$$\frac{d^2 \phi_0}{dt^2} + \omega_1^2 \phi_0 = A \cos((\omega_0 + \omega_1)t - \underline{\kappa \cdot r}) \quad -(17)$$

The usual Euler-Bernoulli structure is:

$$\frac{d^2 \phi_0}{dt^2} + \omega_1^2 \phi_0 = A \cos \omega_0 t, \quad -(18)$$

and at  $\omega_0 = \omega_1$  -(19)

the potential becomes infinite:

$$\phi_0 \rightarrow \infty \quad -(20)$$

It is clear that eq. (18) can be obtained for eq. (17) when:

$$\underline{\kappa \cdot r} = \omega_2 t \quad -(21)$$

where  $\omega_2$  is an angular frequency. Resonance occurs when

$$\omega_0 = \omega_2 \quad -(22)$$

and infinite energy is taken for the system.