

58(5): Spacetime Characteristics of the Whirlpool Galaxy
 The velocity field of fluid spacetime needed for
 formation of a whirlpool galaxy is:

$$\underline{V}_F = \frac{L_F Z}{m(x^2 + y^2)} (-y \underline{i} + x \underline{j}) \quad - (1)$$

So $\underline{\nabla} \cdot \underline{V}_F = 0 \quad - (2)$
 i.e. it is a divergenceless velocity field.

From the Lorenz condition of spacetime:

$$\frac{\partial \Phi_F}{\partial t} + a_0^2 \underline{\nabla} \cdot \underline{V}_F = 0 \quad - (3)$$

it follows that $\frac{\partial \Phi_F}{\partial t} = 0 \quad - (4)$

therefore $\Phi_F = h_F \quad - (5)$

is independent of time, i.e. the spacetime potential or
enthalpy is independent of time
 From the wave equation:

$$\square \Phi_F = \square h_F = 4\pi G \rho_m(\text{matter}) \quad - (6)$$

where $\square = \frac{1}{a_0^2} \frac{\partial^2}{\partial t^2} - \nabla^2 \quad - (7)$

it follows from eq. (4) that:

$$\nabla^2 \Phi_F = -4\pi G \rho_n(\text{matter}) \quad - (8)$$

The gravitational field of the whirlpool galaxy is:

$$\underline{g}(\text{matter}) = - \frac{L_{F2}}{m^2 r_F^3} \underline{e}_r \quad - (9)$$

In spherical polar coordinates:

$$\begin{aligned} \underline{\nabla} \cdot \underline{g} &= \frac{1}{r^2} \frac{d}{dr} (r^2 g_r) \quad - (10) \\ &= \left(\frac{L_{F2}}{m r_F^3} \right)^2 \end{aligned}$$

So the spacetime charge needed for a whirlpool galaxy is:

$$q_F = \left(\frac{L_{F2}}{m r_F^3} \right)^2 \quad - (11)$$

Using the equation:

$$\begin{aligned} \rho_n(\text{matter}) &= \frac{1}{4\pi G} \left(\frac{L_{F2}}{m r_F^3} \right)^2 \quad - (12) \\ &= \frac{q_F}{4\pi G} \end{aligned}$$

gives an expression for the mass density of a whirlpool galaxy.

3) Units Check

The units of G are $m^3 s^{-2} kg^{-1}$ so it is seen that $\rho_m(\text{matter})$ has the correct units of $kg m^{-3}$.

From eq. (12):

$$\rho_m(\text{matter}) \xrightarrow{r_F \rightarrow 0} \infty \quad - (13)$$

So a very large mass density exists at the centre of the galaxy, as observed experimentally.

From eqs. (8) and (12):

$$\nabla^2 \Phi_F = - \left(\frac{L_{F2}}{m r_F^3} \right)^2 \quad - (14)$$

So

$$\Phi_F = h_F = - \frac{1}{6} \left(\frac{L_{F2}}{m} \right)^2 \frac{1}{r_F^3} \quad - (15)$$

The space-time entropy is greatest near the centre of the galaxy.

The force law from eq. (9) is:

$$\begin{aligned} \underline{F}(\text{matter}) &= m \underline{g}(\text{matter}) \\ &= - \frac{L_{F2}}{m r_F^3} \underline{r} \quad - (16) \end{aligned}$$

4) Using the Binet equation:

$$F(r) = -\frac{L^2}{mr^3} \left(\frac{1}{r} + \frac{d^2}{d\theta^2} \left(\frac{1}{r} \right) \right) \quad - (17)$$

The force law (16) corresponds to the hyperbolic spiral orbit:

$$\boxed{\frac{1}{r} = \frac{\theta}{r_0}} \quad - (18)$$

of the star in a whirlpool galaxy.

Therefore the hyperbolic spiral orbit is due to a constant spacetime angular momentum field in fluid gravitation.

As in previous work the orbit (18) leads to the correct velocity curve of the whirlpool galaxy as follows.

In plane polar coordinates (r, θ) , the velocity of a star is:

$$\begin{aligned} v^2 &= \left(\frac{dr}{dt} \right)^2 + r^2 \left(\frac{d\theta}{dt} \right)^2 \quad - (19) \\ &= \left(\frac{d\theta}{dt} \right)^2 \left(r^2 + \left(\frac{dr}{d\theta} \right)^2 \right) \end{aligned}$$

From eq. (18):

$$\left(\frac{dr}{d\theta} \right)^2 = \frac{r^4}{r_0^2} \quad - (20)$$

and for Lagrange's theorem:

$$5) \left(\frac{d\theta}{dt} \right) = \frac{L^2}{m^2 r^4} \quad \text{--- (21)}$$

where L is a constant of motion. So:

$$v^2 = \frac{L^2}{m^2} \left(\frac{1}{r^2} + \frac{1}{r_0^2} \right) \quad \text{--- (22)}$$

$$\xrightarrow{r \rightarrow \infty} \left(\frac{L}{m r_0} \right)^2 = \underline{\text{constant}}$$

as observed experimentally.

As r becomes infinite the velocity of the star becomes constant. In fluid gravitation this is due to a constant spacetime or vacuum or angular momentum.