

1) SUBSIDIARY CONDITION FOR R
Begin with the Evans Lemma:

$$D^\mu (D_\mu q_\nu^a) := 0 \quad - (1)$$

i.e.
$$\partial^\mu (D_\mu q_\nu^a) := 0 \quad - (2)$$

or
$$\partial^\mu \left(\partial_\mu q_\nu^a + \omega_{\mu b}^a q_\nu^b - \Gamma_{\mu\lambda}^\nu q_\nu^a \right) := 0 \quad - (3)$$

$$\Rightarrow \boxed{\square q_\mu^a = R q_\mu^a} \quad - (4)$$

where:

$$R = q_\nu^a \partial^\mu \left(\Gamma_{\mu\lambda}^\nu q_\nu^a - \omega_{\mu b}^a q_\nu^b \right). \quad - (5)$$

Eqn (4) is an eigen equation for R and eqn. (5) is a subsidiary condition for R. In general they must be solved simultaneously. In these notes R is re-expressed in terms only of the connections, the advantage of this is that R is eqn. (4) is an eigenvalue of the tetrad, and R in eqn (5) is a scalar curvature defined

2) similarly to its well known definition in the 1915 theory as a contraction of the Riemann tensor via the Ricci tensor. This eqn (4) is an expression for R in quantum mechanics, and eqn. (5) is an expression for R in classical field theory.

Before embarking on this, proof we note that the tetrad postulate is the same as the first Cartan structure equation. The tetrad postulate is:

$$d_\mu v_\lambda^a + \omega_{\mu b}^a v_\lambda^b = \Gamma_{\mu\lambda}^\sim v_\sim^a \quad - (6)$$

Now use the result:

$$d_\lambda v_\mu^a + \omega_{\lambda b}^a v_\mu^b = \Gamma_{\lambda\mu}^\sim v_\sim^a \quad - (7)$$

and
$$T_{\mu\lambda}^a = T_{\mu\lambda}^\sim v_\sim^a \quad - (8)$$

where the torsion tensor is:

$$T_{\mu\lambda}^\sim = \Gamma_{\mu\lambda}^\sim - \Gamma_{\lambda\mu}^\sim. \quad - (9)$$

Subtracting (7) from (6) we find the first Cartan structure equation:

3)

$$d \wedge q_{\lambda}^a + \omega_{\mu}^a \wedge q_{\lambda}^b = T_{\mu\lambda}^a \quad - (10)$$

or

$$D \wedge q^a = T^a \quad - (10a)$$

Therefore the first Cartan structure equation is the difference of the Ketrand postulates.

Now proceed to develop eqn. (5) using :

$$g^{\mu} = g^{\mu\sigma} \partial_{\sigma} \quad - (11)$$

to give :

$$R = g^{\mu\sigma} q_{\lambda}^a \left(\Gamma_{\mu\lambda}^{\sim} \partial_{\sigma} q_{\sim}^a + q_{\sim}^a \partial_{\sigma} \Gamma_{\mu\lambda}^{\sim} - \omega_{\mu}^a \partial_{\sigma} q_{\lambda}^b - q_{\lambda}^b \partial_{\sigma} \omega_{\mu}^a \right) \quad - (12)$$

Now use the Ketrand postulate again :

$$\partial_{\sigma} q_{\sim}^a + \omega_{\sigma}^a q_{\sim}^b - \Gamma_{\sigma\sim}^{\rho} q_{\rho}^a = 0 \quad - (13)$$

$$\partial_{\sigma} q_{\lambda}^b + \omega_{\sigma}^b q_{\lambda}^c - \Gamma_{\sigma\lambda}^{\sim} q_{\sim}^b = 0 \quad - (14)$$

to find :

4)

$$R = g^{\mu\sigma} \left(g_a^\lambda g_\rho^a \Gamma_{\mu\lambda}^\sim \Gamma_{\sigma\rho}^\sim - g_a^\lambda g_\rho^b \Gamma_{\mu\lambda}^\sim \omega_{\sigma b}^a + g_a^\lambda g_\rho^a \partial_\sigma \Gamma_{\mu\lambda}^\sim - g_a^\lambda g_\rho^b \omega_{\mu b}^a \Gamma_{\sigma\lambda}^\sim + g_a^\lambda g_\rho^c \omega_{\mu b}^a \omega_{\sigma c}^b - g_a^\lambda g_\rho^b \partial_\sigma \omega_{\mu b}^a \right) - (15)$$

Now eliminate ∂ terms using:

$$R^\sigma_{\lambda\sim\mu} = g_a^\sigma g_\lambda^b R^{a b\sim\mu} - (16)$$

$$R^{a b\sim\mu} = \partial_\sim \omega_{\mu b}^a - \partial_\mu \omega_{\sim b}^a + \omega_{\sim c}^a \omega_{\mu b}^c - \omega_{\mu c}^a \omega_{\sim b}^c - (17)$$

$$R^\sigma_{\lambda\sim\mu} = \partial_\sim \Gamma_{\mu\lambda}^\sigma - \partial_\mu \Gamma_{\sim\lambda}^\sigma + \Gamma_{\sim\rho}^\sigma \Gamma_{\mu\lambda}^\rho - \Gamma_{\mu\rho}^\sigma \Gamma_{\sim\lambda}^\rho - (18)$$

$$R^{a b\sim\mu} = g_a^\sigma g_\lambda^b R^\sigma_{\lambda\sim\mu} - (19)$$

The Riemann tensor and ω Riemann form are antisymmetric respectively in σ and λ and in a and b . Using this antisymmetry eqn. (15) reduces to:

$$R = -g^{\mu\sigma} g_a^\lambda g_\rho^b \left(\Gamma_{\mu\lambda}^\sim \omega_{\sigma b}^a + \omega_{\mu b}^a \Gamma_{\sigma\lambda}^\sim \right) - (20)$$

5) Now simplify and remove Φ terms using:

$$g_{\nu}^{\lambda} \Gamma_{\mu\lambda}^{\sim} = \Gamma_{\mu a}^{\sim} \quad - (21)$$

$$g_{\nu}^b \omega_{\sigma b}^a = \omega_{\sigma \nu}^a \quad - (22)$$

$$g_{\nu}^{\lambda} \omega_{\mu b}^a = \omega_{\mu b}^{\lambda} \quad - (23)$$

$$g_{\nu}^b \Gamma_{\sigma\lambda}^{\sim} = \Gamma_{\sigma\lambda}^b \quad - (24)$$

to give:

$$R = -g^{\mu\sigma} \left(\Gamma_{\mu a}^{\sim} \omega_{\sigma \nu}^a + \omega_{\mu b}^{\lambda} \Gamma_{\sigma\lambda}^b \right) \quad - (25)$$

Finally re-arrange dummy indices $b \rightarrow a$; $\lambda \rightarrow \nu$:

$$R = -g^{\mu\sigma} \left(\Gamma_{\mu a}^{\sim} \omega_{\sigma \nu}^a + \omega_{\mu a}^{\nu} \Gamma_{\sigma \nu}^a \right) \quad - (26)$$

Eqn (26) is the desired generalization of the equation for R in the 1915 theory:

$$R = g^{\mu\sigma} R_{\mu\sigma} \quad - (27)$$

6) (comparing eqns (26) and (27) gives Φ Ricci tensor is Φ Evans unified field theory:

$$R_{\mu\nu} = - \left(\Gamma_{\mu a}^{\sim} \omega_{\sigma\nu}^a + \omega_{\mu a}^{\sim} \Gamma_{\sigma\nu}^a \right) \quad - (28)$$

Using Φ index-contracted Euler field eq:

$$R = -kT \quad - (29)$$

it is found that:

$$kT = g^{\mu\nu} \left(\Gamma_{\mu a}^{\sim} \omega_{\sigma\nu}^a + \omega_{\mu a}^{\sim} \Gamma_{\sigma\nu}^a \right) \quad - (30)$$

Eqn. (30) is Φ classical subsidiary condition of Φ Evans wave equation of quantum mechanics:

$$\left(\square + kT \right) \psi_{\mu}^a = 0. \quad - (31)$$