

327(8): Precession from the Minkowski Metric

As in previous notes the orbit from the Minkowski metric of special relativity is:

$$r = \frac{d}{1 + \epsilon \cos \theta} \quad - (1)$$

where d and ϵ are the respectively the relativistic half right latitude and the relativistic ellipticity. Therefore:

$$\cos \theta = \frac{1}{\epsilon} \left(\frac{d}{r} - 1 \right) \quad - (2)$$

Define θ by: $\theta = \theta_0 + \Delta\theta \quad - (3)$

where $\Delta\theta$ is the precession angle of the Minkowski metric.

At the perihelion:

$$\theta_0 = 2\pi \quad - (3)$$

and

$$r = r_{\min} = \frac{d_0}{1 + \epsilon_0} \quad - (4)$$

of the classical or Newtonian orbit with:

$$r = \frac{d_0}{1 + \epsilon_0 \cos \theta_0} \quad - (5)$$

Therefore at the perihelion:

$$\cos(2\pi + \Delta\theta) = \cos \Delta\theta \quad - (6)$$

and $\sin \Delta\theta \sim \Delta\theta = (1 - \cos^2 \Delta\theta)^{1/2} \quad - (7)$

It follows from eqs. (2) and (4) that:

$$\cos \Delta\theta = \frac{1}{\epsilon} \left(\frac{d}{d_0} (1 + \epsilon_0) - 1 \right) - (8)$$

At the perihelion. As :

$$d \rightarrow d_0 \text{ and } \epsilon \rightarrow \epsilon_0 - (9)$$

eq. (8) gives :

$$\Delta\theta \rightarrow 0 - (10)$$

self consistently.

Therefore the precession is given by :

$$\Delta\theta = \left(1 - \frac{1}{\epsilon^2} \left(\frac{d}{d_0} (1 + \epsilon_0) - 1 \right)^2 \right)^{1/2} - (11)$$

where

$$\frac{d}{d_0} = \gamma^2 = \left(1 - \left(\frac{v_0}{c} \right)^2 \right)^{-1} - (12)$$

The relativistic ϵ can be worked out as in previous notes, and to an excellent approximation

$$\epsilon \sim \epsilon_0 - (13)$$

Therefore :

$$\Delta\theta = \left(1 - \frac{1}{\epsilon_0^2} \left(\frac{d}{d_0} (1 + \epsilon_0) - 1 \right)^2 \right)^{1/2} - (14)$$
