

262(1): The Development of a Force Law for Any Orbit

Consider the orbital and spin force laws given in UFT

261:

$$\underline{F}(\text{orb}) = -\underline{\nabla} \phi - \frac{\partial \underline{p}}{\partial t} - \underline{\omega}_0 \underline{p} + \underline{\phi} \underline{\omega} \quad - (1)$$

$$\underline{F}(\text{spin}) = \underline{\nabla} \times \underline{p} - \underline{\omega} \times \underline{p} \quad - (2)$$

By antisymmetry:

$$-\frac{\partial \underline{p}}{\partial t} - \underline{\omega}_0 \underline{p} = -\underline{\nabla} \phi + \underline{\omega} \phi \quad - (3)$$

Eq. (3) is obtained from the definition of torque in

Cartan geometry:

$$T^a_{\mu\nu} = \partial_\mu q^a_\nu - \partial_\nu q^a_\mu + \omega^a_{\mu b} q^b_\nu - \omega^a_{\nu b} q^b_\mu \quad - (4)$$

where:

$$\Gamma^a_{\mu\nu} = \partial_\mu q^a_\nu + \omega^a_{\mu b} q^b_\nu \quad - (5)$$

$$\Gamma^a_{\nu\mu} = \partial_\nu q^a_\mu + \omega^a_{\nu b} q^b_\mu \quad - (6)$$

By the commutator theorem of geometry:

$$\Gamma^a_{\mu\nu} = -\Gamma^a_{\nu\mu} \quad - (7)$$

So:

$$2) \partial_\mu q_\nu^a + \omega_{\mu b}^a q_\nu^b = -(\partial_\mu q_\nu^a + \omega_{\mu b}^a q_\nu^b) - (8)$$

For example:

$$\partial_0 q_1^a + \omega_{0b}^a q_1^b = -(\partial_1 q_0^a + \omega_{1b}^a q_0^b) - (9)$$

In vector notation, eq. (9) is:

$$-\frac{1}{c} \frac{d \underline{q}}{dt} - \underline{\omega}_{0b}^a \underline{q}^b = -\underline{\nabla} q_0^a + \underline{\omega}^a_b q_0^b - (10)$$

With the definitions:

$$\omega_{\mu b}^a = \left(\frac{1}{c} \omega_{0b}^a, -\underline{\omega}^a_b \right) - (11)$$

and $\phi^a_0 = c q_0^a, - (12)$

$$p_\mu^a = \left(\frac{\phi^a_0}{c}, -\underline{p}^a \right) - (13)$$

and using the polarization model for simplicity only:

$$-\frac{d \underline{p}}{dt} - \underline{\omega}_0 \underline{p} = -\underline{\nabla} \phi + \underline{\omega} \phi - (14)$$

In the absence of a spin current eq. (14) becomes:

$$-\frac{d\phi}{dt} = -\nabla \phi \quad (15)$$

which is the equivalence of gravitational and inertial mass, tested experimentally to one part in ten power seventeen. Therefore eq. (14) is the generally covariant format of the equivalence theorem.

It is convenient to define the force as:

$$\underline{F} = -\frac{d\phi}{dt} - \omega \cdot \phi = -\nabla \phi + \underline{\omega} \phi \quad (16)$$

and to compare this definition with the equivalent definition from the plane polar representation of any planar orbit, in which the velocity is:

$$\underline{v} = \dot{r} \underline{e}_r + r \dot{\theta} \underline{e}_\theta \quad (17)$$

and the acceleration is:

$$\underline{a} = (\ddot{r} - r \dot{\theta}^2) \underline{e}_r + (r \ddot{\theta} + 2 \dot{r} \dot{\theta}) \underline{e}_\theta \quad (18)$$

As argued in previous work eqs. (17) and (18) are examples of Curvilinear geometry because the axes of the plane polar system are rotating. So there is present a geometrical

4) connection and eqs. (17) and (18) are equations of general relativity for \mathcal{O} , reason. Eq. (16) is also an equation of general relativity. If the axes were not rotating then:

$$\underline{v} = \dot{r} \underline{e}_r \quad - (19)$$

$$\underline{a} = \ddot{r} \underline{e}_r \quad - (20)$$

and this is the Newtonian theory.

For any orbit is a plane, a Lagrangian analysis shows that the angular momentum is a constant of motion and is conserved. It is given by:

$$L = m r^2 \dot{\theta} \quad - (21)$$

Therefore:
$$\dot{\theta} = \frac{L}{m r^2} \quad - (22)$$

and
$$\dot{r} = \frac{dr}{d\theta} \frac{d\theta}{dt} = \dot{\theta} \frac{dr}{d\theta} \quad - (23)$$

Let
$$f(r) = \frac{L}{m r^2} \quad - (24)$$

then
$$\ddot{\theta} = \frac{df(r)}{dt} = \frac{df(r)}{dr} \frac{dr}{dt} \quad - (25)$$

So:

$$\ddot{\theta} = \dot{r} \frac{d}{dr} \left(\frac{L}{mr^2} \right) = -\frac{2L}{mr^3} \dot{r} \quad (26)$$

So

$$r \ddot{\theta} = -\frac{2L}{mr^2} \dot{r} \quad (27)$$

and

$$2 \dot{r} \dot{\theta} = \frac{2L}{mr^2} \dot{r} \quad (28)$$

It follows that the Coriolis force for any planar orbit vanishes:

$$\begin{aligned} \underline{F}_{\text{Coriolis}} &= m (r \ddot{\theta} + 2 \dot{r} \dot{\theta}) \underline{e}_{\theta} \\ &= \underline{0} \end{aligned} \quad (29)$$

The force from eq. (18) for any planar orbit is therefore radial:

$$\begin{aligned} \underline{F} &= m \underline{a} = (\ddot{r} - r \dot{\theta}^2) \underline{e}_r \\ &= \left(\frac{d^2 r}{dt^2} - \frac{L^2}{m r^3} \right) \underline{e}_r \end{aligned} \quad (30)$$

The Lagrangian analysis produces

b) the equation:

$$F(r) = -\frac{L}{mr^2} \left(\frac{d^2}{dt^2} \left(\frac{1}{r} \right) + \frac{1}{r} \right) \quad (31)$$

and this is again an equation of general relativity because it uses the rotating axes of the plane polar coordinate system.

If the orbit is a static ellipse then:

$$\frac{1}{r} = \frac{1}{a} (1 + e \cos \theta) \quad (32)$$

so

$$\frac{d^2}{dt^2} \left(\frac{1}{r} \right) = -\frac{e}{a} \cos \theta \quad (33)$$

and

$$F(r) = -\frac{L}{mr^2 a} \quad (34)$$

The orbit has no Coriolis force so:

$$\underline{F}(r) = -\frac{L}{mdr^2} \underline{e}_r \quad (35)$$

From eqs (30) and (35):

$$\frac{d^2 r}{dt^2} - \frac{L^2}{mr^3} = -\frac{L}{mdr^2} \quad (36)$$

so:

$$\frac{d^2 r}{dt^2} = -\frac{L}{m d r^2} + \frac{L^2}{m r^3} \quad - (37)$$

In a stable orbit:

$$\frac{d^2 r}{dt^2} = 0 \quad - (38)$$

so the inverse square law of attraction is balanced by the outward force proportional to $1/r^3$. This is the centrifugal force:

$$\begin{aligned} -\underline{\omega} \times (\underline{\omega} \times \underline{r}) m &= \frac{L^2}{m r^3} \underline{e}_r \quad - (39) \\ &= r \dot{\theta}^2 \underline{e}_r. \end{aligned}$$

From vector analysis:

$$\underline{\omega} \times (\underline{\omega} \times \underline{r}) = \underline{\omega} (\underline{\omega} \cdot \underline{r}) - \omega^2 \underline{r}. \quad - (40)$$

Since

$$\underline{\omega} = \frac{d\theta}{dt} \underline{k} \quad - (41)$$

$$-\underline{\omega} \times (\underline{\omega} \times \underline{r}) = \omega^2 r \underline{e}_r \quad - (42)$$

Eqn. (39), e.c.d.

8) The centrifugal force is a result of the rotation of the axes in a plane polar coordinate system. It does not occur in Newtonian dynamics and is an example of Cartan geometry, which always governs physics. It is obviously a reasonable force.

From eqs (14) and (30):

$$\begin{aligned}\underline{F} &= -\frac{d\underline{p}}{dt} - \underline{\omega} \cdot \underline{p} = -\underline{\nabla} \phi + \underline{\omega} \phi \\ &= \left(\frac{d^2 r}{dt^2} - \frac{L^2}{m r^3} \right) \underline{e}_r - (43)\end{aligned}$$

and

$$\underline{F} = -\frac{L}{mr^2} \left(\frac{d^2}{dt^2} \left(\frac{1}{r} \right) + \frac{1}{r} \right) \underline{e}_r - (44)$$

because \underline{e}_θ is no planar dependent Coriolis force in any orbit, so \underline{F} is always radial.

These equations are valid for any lower orbit including the orbit of a star in

a whirlpool galaxy, and for all perihelia precessions.
 The Newtonian result is:.

$$\underline{F} = -\underline{\nabla} \phi = \frac{d^2 \underline{r}}{dt^2} \underline{e}_r - (45)$$

from eq. (43), and:

$$\underline{F} = -\frac{L}{m d r^2} \underline{e}_r - (46)$$

from eq. (44) for an elliptical orbit. So the
Newtonian dynamics the centrifugal term is
missing.

The centrifugal term is due to the spin
 connection, and precession is UFT261:

$$\phi \underline{\omega} = -\frac{L^2}{m r^3} \underline{e}_r - (47)$$

In a near Newtonian approximation:

$$\phi = -\frac{m M G}{r} - (48)$$

here m is attracted to M at the focus
 of an ellipse and G is the Newton constant.
 So :

$$-\frac{mM G}{r} \underline{\omega} = -\frac{L^2}{m r^3} \underline{e}_r - (49)$$

and the spi connection is :

$$\underline{\omega} = \frac{L^2}{m^2 M G r^2} \underline{e}_r - (50)$$

The right latitude d in Newton theory is :

$$d = \frac{L^2}{m^2 M G} - (51)$$

so

$$\underline{\omega} = \frac{d}{r^2} \underline{e}_r - (52)$$

Eq. (50) is true for all ^{planar} orbits, and eq. (52) for an ellipse.
