

132(9) : Equivalence of Electrodynamics Metric and Minimal Prescription

In spherically symmetric spacetime the electrodynamic metric is:

$$ds^2 = c^2 d\tau^2 = e^{-r_0/r} c^2 dt^2 - e^{r_0/r} dr^2 - r^2 d\phi^2 \quad - (1)$$

in the xy plane is cylindrical polar coordinates. Here:

$$r_0 = 2 \frac{e_1 e_2}{m \frac{4\pi\epsilon_0}{c^2}} \quad - (2)$$

If L is the Lagrangian, H the Hamiltonian, and T the kinetic energy, then:

$$L = H = T = \frac{1}{2} mc^2 = \frac{1}{2} m \left(e^{-r_0/r} c^2 \left(\frac{dt}{d\tau} \right)^2 - e^{r_0/r} \left(\frac{dr}{d\tau} \right)^2 - r^2 \left(\frac{d\phi}{d\tau} \right)^2 \right) \quad - (3)$$

The Lagrange equation is:

$$\frac{d}{d\tau} \left(\frac{\partial L}{\partial \dot{x}^\mu} \right) = \frac{\partial L}{\partial x^\mu} \quad - (4)$$

Therefore the constants of motion are the conserved quantities:

$$E = mc^2 e^{-r_0/r} \frac{dt}{d\tau} = \text{constant} \quad - (5)$$

$$p = m e^{r_0/r} \frac{dr}{d\tau} = \text{constant} \quad - (6)$$

$$L = m r^2 \frac{d\phi}{d\tau} = \text{constant} \quad - (7)$$

These are the total energy (E), linear momentum (p) and angular momentum (L). So:

$$T = \frac{1}{2} mc^2 = \frac{1}{2} \left(e^{r_0/r} \frac{E^2}{mc^2} - e^{-r_0/r} \frac{p^2}{m} - \frac{L^2}{mr^2} \right) \quad - (8)$$

2) I_2 limit: $\frac{r_0}{r} \rightarrow 0$ — (9)

$$H = \frac{1}{2} mc^2 = \frac{1}{2} \left(\frac{E^2}{mc^2} - \frac{p^2}{m} - \frac{L^2}{mr^2} \right) \text{ — (10)}$$

If there is no angular momentum:
 $L = 0$ — (11)

then $H = \frac{1}{2} mc^2 = \frac{1}{2} \left(\frac{E^2}{mc^2} - \frac{p^2}{m} \right)$ — (12)

i.e. $E^2 = c^2 p^2 + m^2 c^4$ — (13)

which is the Einstein energy equation of special relativity. The Hamiltonian is:

$$H = \frac{1}{2m} p^\mu p_\mu \text{ — (14)}$$

where

$$p^\mu = \left(\frac{E}{c}, \underline{p} \right) \text{ — (15)}$$

$$p_\mu = \left(\frac{E}{c}, -\underline{p} \right) \text{ — (16)}$$

— (17)

For finite r_0/r eqn. (12) is:

$$H = \frac{1}{2m} \left(\left(\exp\left(\frac{r_0}{2r}\right) \frac{E}{c} \right)^2 - \left(\exp\left(-\frac{r_0}{2r}\right) \underline{p} \right)^2 \right)$$

Consider the case:

$$\exp\left(\frac{r_0}{2r}\right) \sim 1 + \frac{r_0}{2r} \text{ — (18)}$$

$$\exp\left(-\frac{r_0}{2r}\right) \sim 1 - \frac{r_0}{2r} \text{ — (19)}$$

Then the change in the Hamiltonian due to $r_0/(2r)$ can be represented by:

$$E \rightarrow E + \left(\frac{r_0}{2r}\right)E - (20)$$

$$P \rightarrow P - \left(\frac{r_0}{2r}\right)P - (21)$$

This representation is a minimal prescription:

$$p^\mu \rightarrow p^\mu - e_1 A^\mu - (22)$$

where

$$A^\mu = \left(\frac{\phi}{c}, \underline{A}\right) - (23)$$

if

$$\frac{r_0 E}{2r} = -e_1 \phi - (24)$$

$$\frac{r_0 P}{2r} = e_1 A - (25)$$

i.e.

$$\boxed{\begin{aligned} \phi &= -\frac{e_2}{4\pi\epsilon_0 r} \left(\frac{E}{mc^2}\right) \\ A &= \frac{e_2}{4\pi\epsilon_0 r c} \left(\frac{P}{mc}\right) \end{aligned}} - (26)$$

For a particle at rest:

$$\phi = -\frac{e_2}{4\pi\epsilon_0 r} - (27)$$

$$A = 0 - (28)$$

and self-consistently, the problem is one of electrodynamics.
 If the definition (26) the metric gives the relativistic
Milner-Jacobi equation.

$$(p^\mu - eA^\mu)(p_\mu - eA_\mu) = m^2 c^2 - (29)$$

and the correct Dirac equation of an electron is a

4) potential A^μ :

$$(i \gamma^\mu (\partial_\mu - \frac{ieA_\mu}{\hbar}) - \frac{mc}{\hbar}) \psi = 0 \quad (30)$$

$$\text{or } (\gamma^\mu (\partial_\mu - eA_\mu) - mc) \psi = 0 \quad (31)$$

This gives the g factor of the electron, the precessional effect of Thomas is spin orbit coupling, and ESR and NMR.

Conclusion The metric (1) correctly reduces to the Dirac equation (31). For electrostatics:

$$A_\mu = \left(\frac{\phi}{c}, 0 \right) \quad (32)$$

In this case, using well known methods:

$$H \doteq \frac{p^2}{2m} + e\phi \quad (33)$$

which is the non-relativistic limit of the RHS of eq. (26) of note 152(8).

The most general Hamiltonian is eq. (8), and this generalizes the Dirac equation (31), giving new spectral effects.
