

**SOME NOTES ON THE SOLENOIDAL BELTRAMI EQUATIONS**  
(Barrett and Grimes, pp. 228 ff).

This is

$$\begin{aligned}\nabla \times \mathbf{B} &= k\mathbf{B} \\ \nabla \cdot \mathbf{B} &= 0\end{aligned}$$

Note that this is satisfied by:

$$\begin{aligned}\mathbf{B}^{(1)} &= \frac{B^{(0)}}{\sqrt{2}}(i\mathbf{i} + j\mathbf{j})e^{i(\omega t - \kappa \cdot \mathbf{r})} \\ \mathbf{B}^{(2)} &= \frac{B^{(0)}}{\sqrt{2}}(-i\mathbf{i} + j\mathbf{j})e^{-i(\omega t - \kappa \cdot \mathbf{r})} \\ \mathbf{B}^{(3)} &= B^{(0)}\mathbf{k}\end{aligned}$$

Specifically:

$$\begin{aligned}\nabla \times \mathbf{B}^{(1)} &= \begin{bmatrix} i & j & k \\ \frac{\partial}{\partial X} & \frac{\partial}{\partial Y} & \frac{\partial}{\partial Z} \\ ie^{\phi} & e^{\phi} & 0 \end{bmatrix} \frac{B^{(0)}}{\sqrt{2}} \\ &= \left( -i \frac{\partial}{\partial Z} e^{\phi} + ij \frac{\partial}{\partial Z} e^{\phi} \right) \frac{B^{(0)}}{\sqrt{2}} \\ &= \kappa (i\mathbf{i} + j\mathbf{j}) e^{\phi} \frac{B^{(0)}}{\sqrt{2}} \\ &= \kappa \mathbf{B}^{(1)}\end{aligned}$$

$\nabla \times \mathbf{B}^{(1)} = \kappa \mathbf{B}^{(1)}; \quad \nabla \cdot \mathbf{B}^{(1)} = 0$	(1)
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Similarly:

$$\begin{aligned}
 \nabla \times \mathbf{B}^{(2)} &= \begin{bmatrix} i & j & k \\ \frac{\partial}{\partial X} & \frac{\partial}{\partial Y} & \frac{\partial}{\partial Z} \\ -ie^{-\phi} & e^{-\phi} & 0 \end{bmatrix} \frac{B^{(0)}}{\sqrt{2}} \\
 &= \left( -i \frac{\partial}{\partial Z} e^{-\phi} - ij \frac{\partial}{\partial Z} e^{-\phi} \right) \frac{B^{(0)}}{\sqrt{2}} \\
 &= \kappa (-ii + j) e^{-\phi} \frac{B^{(0)}}{\sqrt{2}} \\
 &= \kappa \mathbf{B}^{(2)}
 \end{aligned}$$

$$\nabla \times \mathbf{B}^{(2)} = \kappa \mathbf{B}^{(2)}; \quad \nabla \cdot \mathbf{B}^{(2)} = 0 \quad (2)$$

Similarly:

$$\nabla \times \mathbf{B}^{(3)} = \mathbf{0}; \quad \nabla \cdot \mathbf{B}^{(3)} = 0. \quad (3)$$

## CONCLUSION

The field  $\mathbf{B}^{(1)}$ ,  $\mathbf{B}^{(2)}$ ,  $\mathbf{B}^{(3)}$  are solenoidal and Beltrami, and it is clear that this does not mean that  $\mathbf{B}^{(1)}$ ,  $\mathbf{B}^{(2)}$  and  $\mathbf{B}^{(3)}$  are everywhere constant.

### THE CORRECT INTERPRETATION OF $\nabla \times B^{(3)}$ AND $\nabla \cdot B^{(3)}$

This is:

$$\nabla \times B^{(3)} = -ig \nabla \times (A^{(1)} \times A^{(2)})$$

$$\nabla \cdot B^{(3)} = -ig \nabla \cdot (A^{(1)} \times A^{(2)})$$

which shows that  $B^{(3)}$  is not everywhere constant, because  $A^{(1)} \times A^{(2)}$  is not everywhere constant. Specifically:

1)

$$\begin{aligned} \nabla \cdot B^{(3)} &= -ig \left[ A^{(2)} \cdot (\nabla \times A^{(1)}) - A^{(1)} \cdot (\nabla \times A^{(2)}) \right] \\ &= -ig (A^{(2)} \cdot B^{(1)} - A^{(1)} \cdot B^{(2)}) \\ &= 0 \end{aligned}$$

It is clear that  $A^{(2)} = A^{(1)*}$  are defined as propagating fields, and so  $B^{(3)}$  is a propagating field that exists if and only if  $A^{(1)} \times A^{(2)}$  exists, at a phase number  $\omega t - \kappa \cdot Z$ .

2)

$$\begin{aligned} \nabla \times B^{(3)} &= ig \left[ A^{(1)} (\nabla \cdot A^{(2)}) - (\nabla \cdot A^{(1)}) A^{(2)} + (A^{(2)} \cdot \nabla) A^{(1)} - (A^{(1)} \cdot \nabla) A^{(2)} \right] \\ &= 0 \end{aligned}$$

3)

$$\begin{aligned} \frac{\partial B^{(3)}}{\partial t} &= -ig \frac{\partial}{\partial t} (A^{(1)} \times A^{(2)}) \\ &= -ig \left( \frac{\partial A^{(1)}}{\partial t} \times A^{(2)} + A^{(1)} \times \frac{\partial A^{(2)}}{\partial t} \right) \\ &= 0 \end{aligned}$$