

## AN EXPLANATION OF SNELL'S LAW USING NON-ABELIAN ELECTRODYNAMICS

### ABSTRACT

The conventional explanation for simple reflection and Snell's law in Maxwell-Heaviside theory violates the parity inversion operator  $P$ . It is shown that a self-consistent explanation is given in non-Abelian electrodynamics, using the appropriate Stokes theorem for the electromagnetic phase.

### INTRODUCTION

Snell's Law and normal reflection in optics are of course well verified empirically, but the Maxwell-Heaviside explanation {1} of this simple phenomenon of optics is shown to violate parity inversion symmetry ( $P$ ). It is suggested that a non-Abelian electromagnetic phase, based on the appropriate Stokes theorem {2}, gives a self-consistent description of the effect, and equates the dynamical and topological phases. The topological phase is missing entirely from the Maxwell-Heaviside theory, but is well observed {3-5} in interferometry.

### P VIOLATION

The Maxwell-Heaviside theory of electrodynamics is a U(1) gauge theory {6, 7}, so is referred to, for brevity, simply as "U(1)". Reflection is described in terms of the incident U(1) phase, and the reflected U(1) phase. The former is

$$i(\boldsymbol{\kappa} \cdot \boldsymbol{r} - \omega t) \quad (1)$$

and the latter for normal reflection is

$$i(-\boldsymbol{\kappa} \cdot \boldsymbol{r} - \omega t) \quad (2)$$

where  $\boldsymbol{\kappa}$  is the wave-vector at position  $\boldsymbol{r}$ , and  $\omega$  the frequency at instant  $t$ . Normal reflection is equivalent to parity inversion,  $P$ , whose effect is:

$$P(X, Y, Z) = -X, -Y, -Z \quad (3)$$

if  $\boldsymbol{r}$  is in the  $Z$  axis then  $P(Z) = -Z$  by definition, and  $P(\boldsymbol{r}) = -\boldsymbol{r}$ . However,  $P(\boldsymbol{\kappa}) = -\boldsymbol{\kappa}$  and  $P(\omega t) = \omega t$ . Therefore:

$$P(\boldsymbol{\kappa} \cdot \boldsymbol{r} - \omega t) = \boldsymbol{\kappa} \cdot \boldsymbol{r} - \omega t \quad (4)$$

and the usual U(1) description of normal reflection violates parity inversion symmetry and is invalid. This is easily seen from the fact that the U(1) phase is a number invariant under  $P$ , and motion reversal  $T$ . Therefore U(1) is unable to describe Sagnac and Michelson interferometry, as has been recently realized {8-11}.

## SELF-CONSISTENT DESCRIPTION OF SNELL'S LAW IN NON-ABELIAN ELECTRODYNAMICS.

If electrodynamics is described by an  $O(3)$  gauge theory, {12-15} the Snell Law originates in a phase which is described by a non-Abelian Stokes theorem. {16} It has been shown elsewhere {10} that this can be reduced to:

$$\oint \kappa \cdot dZ = g \int B^{(3)} \cdot dS \quad (5)$$

where the left hand side is a line integral, and the right hand side is an area integral over the Evans-Vigier field  $B^{(3)}$ {8-11}. If a beam of light originates at  $O$  and is normally reflected from a perfectly reflecting mirror at the point  $Z$ , the line integral is:

$$\oint \kappa \cdot dZ = \int_0^Z \kappa dZ - \int_Z^0 \kappa dZ = 2\kappa Z. \quad (6)$$

Notice that this gives, by chance, the same phase change,  $2\kappa Z$  as in the  $U(1)$  description of normal reflection, which is therefore fortuitously useful as a calculating device, but physically incorrect.

The area integral on the right hand side of eqn. (5) is a topological phase {3-5}. Using the definition  $g = \kappa A^{(0)}$  and

$$B^{(3)} = \kappa A^{(0)} \quad (7)$$

the right hand side becomes  $\kappa^2 S$ , where  $S$  is an area:

$$S = \frac{2Z}{\kappa}. \quad (8)$$

If the distance  $OZ$  is  $n$  wavelengths,  $\lambda$ , then the area is:

$$S = \frac{n\lambda^2}{\pi} \quad (9)$$

Normal reflection is a special case of the Snell Law, and the latter can always be reduced to normal reflection by utilizing the components in  $Z$  of the incident and reflected beams. These components are equal and opposite as in eqn. (6) because the incident and reflected angles to the normal are equal by Snell's Law.

## DISCUSSION

The outcome of this very simple example is that all of electrodynamics (classical and quantum) must be upgraded to a gauge theory of higher symmetry than  $U(1)$ . The  $O(3)$  gauge symmetry proves to be effective {8-11}. The Snell Law is of course empirically well verified, but the  $U(1)$  description of it violates  $P$ . The  $O(3)$  description in eqn. (5) is self-consistent, because under  $P$ , both sides are negative. The left hand side is negative because the line integral changes sign under  $P$ , and the right hand side is negative because the integral is negative under  $P$  (product of an axial vector  $B^{(3)}$  and a polar vector  $S$ ). A rigorous derivation of eqn. (5) is available elsewhere {10}.

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- {16} A transparent form of this theorem is given by Broda in ref. (2).